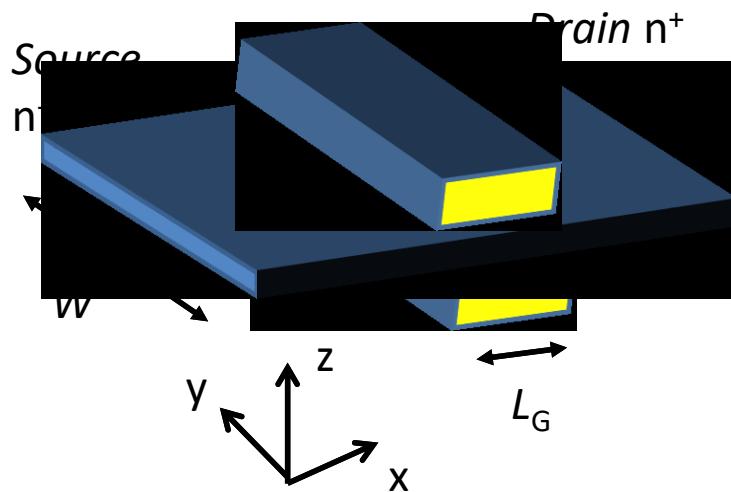


Lecture 3&4 –Ballistic 2D FETs

- 2D Ballistic FET (83-114)
- General Expression for currents
 - 2D effects
 - Bias self-consistency

2D Ballistic FET



Carriers are confined to the xy-plane

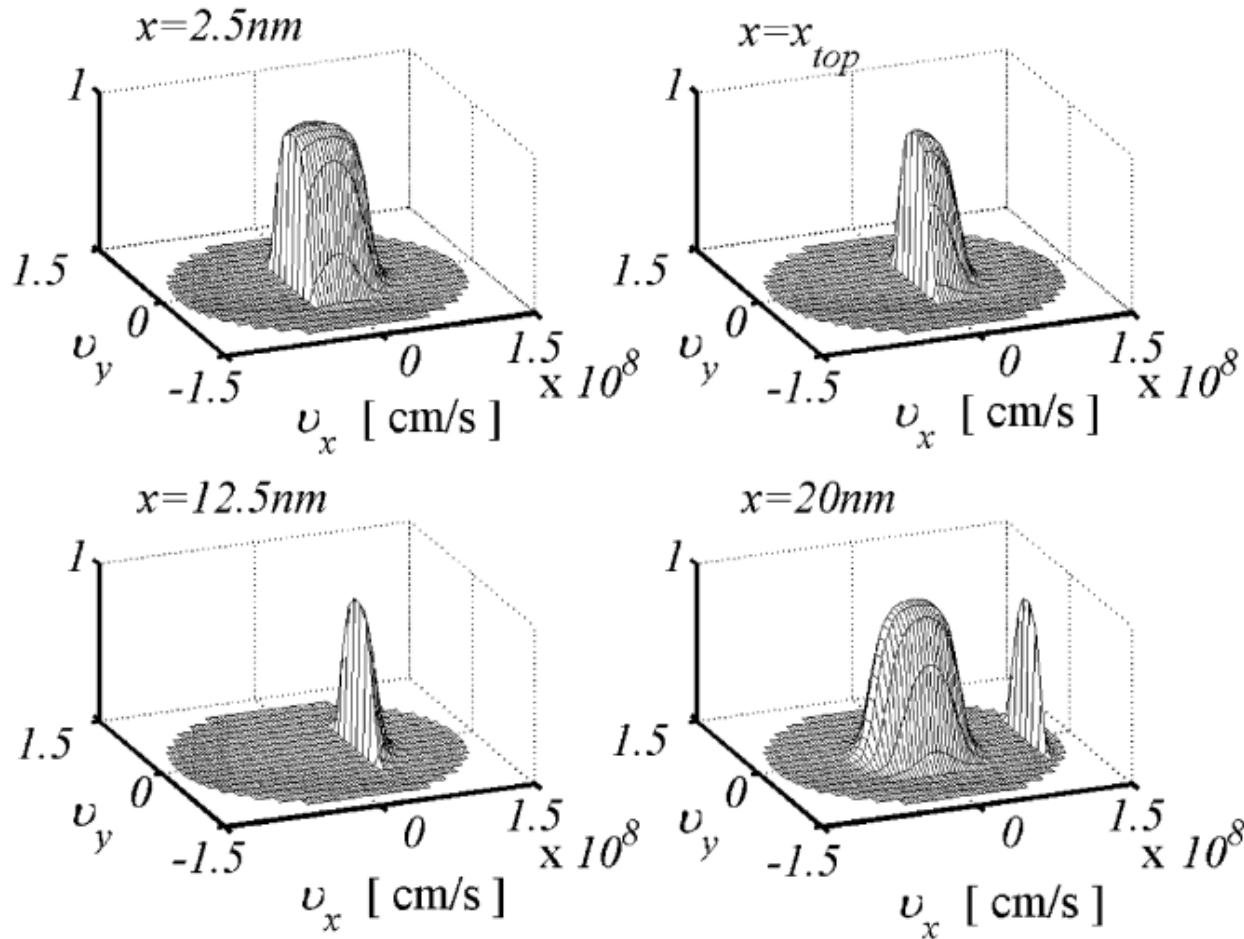
Electron velocities – distributed as v_x and v_y
 v_x gives conduction current

Subbands $E_1(x)$, $E_2(x)$

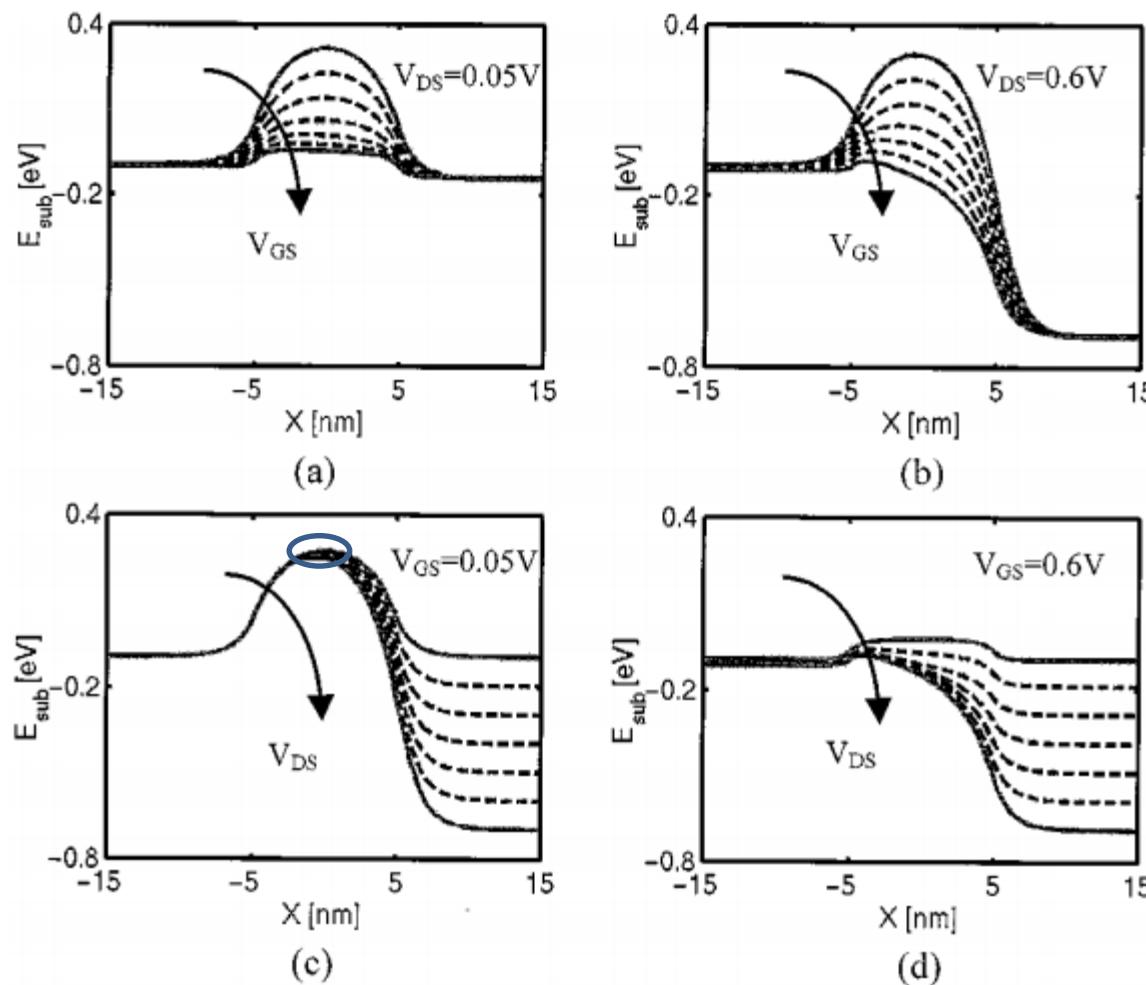
The device width W is assumed to be large –
neglect quantization in y direction

Numerical 2D Ballistic Modeling

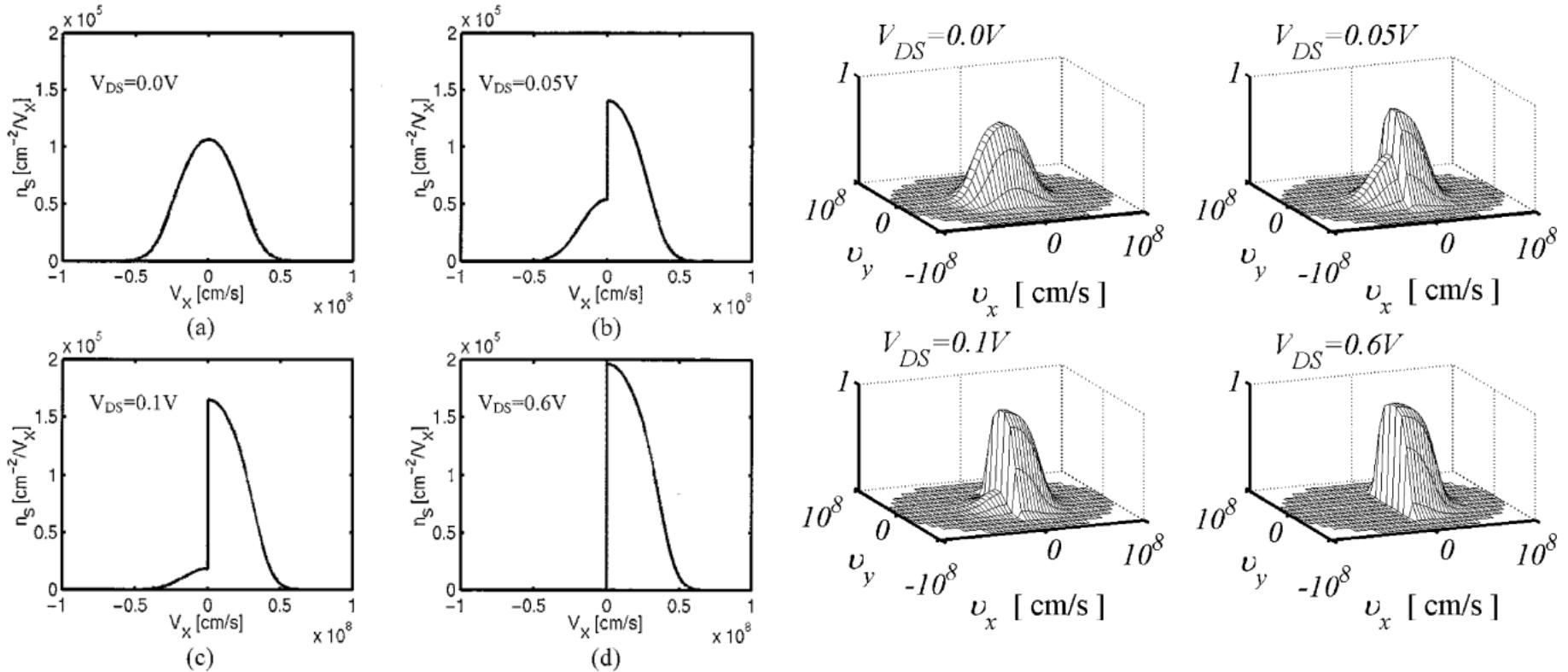
J.-H. Rhew et al. / Solid-State Electronics 46 (2002) 1899–1906



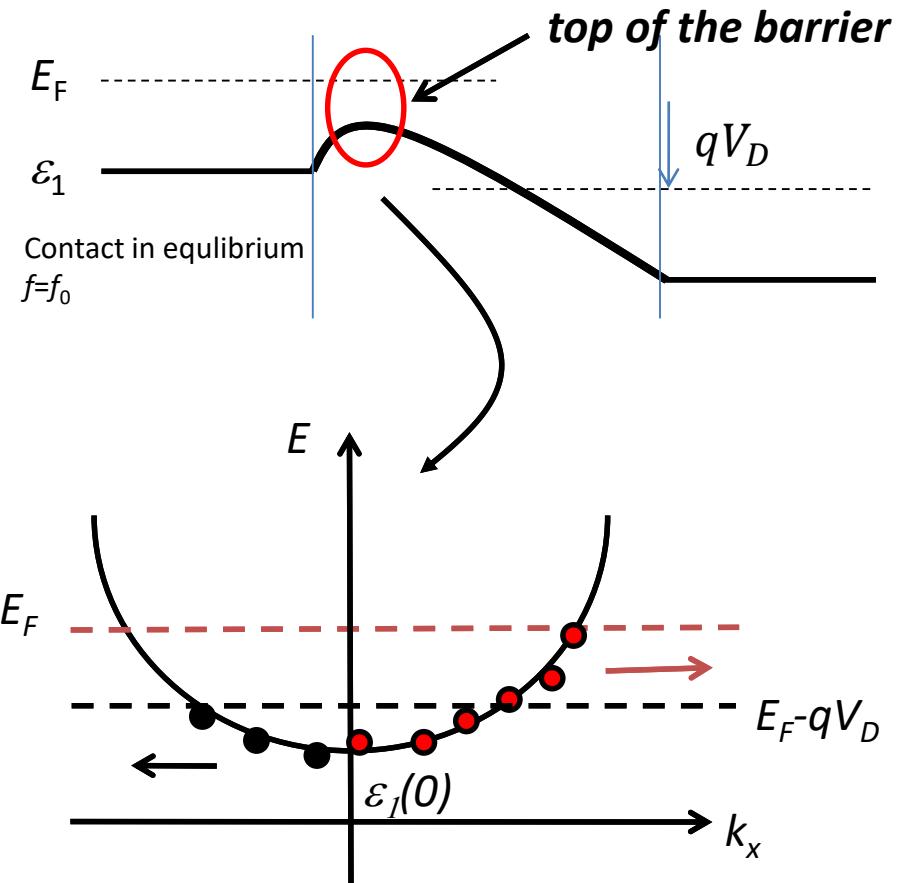
Numerical 2D Ballistic Modeling



Current Saturation



2D Ballistic MOSFET Currents / Charges



$$n_s^+(0) = \frac{1}{A} \sum_{k_x, k_y > 0} f_0(E_F)$$

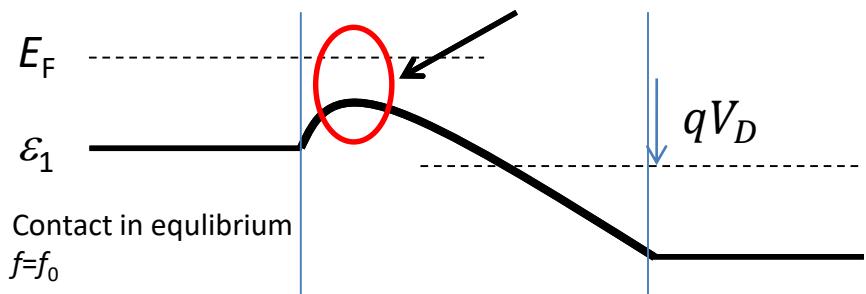
$$J^+ = \frac{1}{A} \sum_{k_x, k_y > 0} q v_x f_0(E_F)$$

$$n_s^-(0) = \frac{1}{A} \sum_{k_x, k_y < 0} f_0(E_F - qV_D)$$

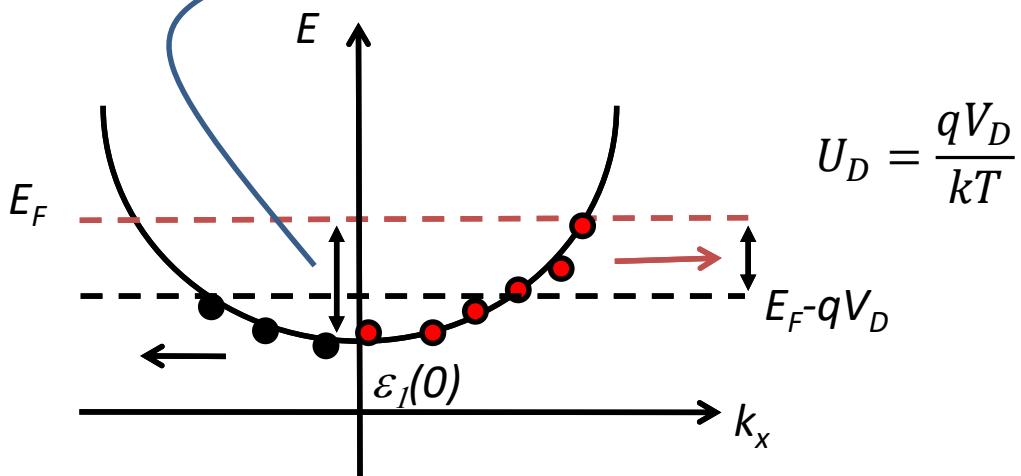
$$J^- = \frac{1}{A} \sum_{k_x, k_y < 0} q v_x f_0(E_F - qV_D)$$

$$\sum_{\mathbf{k}} g(\mathbf{k}) \rightarrow \frac{A}{(2\pi)^2} \int_{\mathbf{k}} g(\mathbf{k}) d\mathbf{k}$$

2D Ballistic MOSFET Currents / Charges



$$\eta_F = (E_F - \varepsilon_1)/kT$$



$$n_s^+(0) = \frac{N_{2D}}{2} F_0(\eta_F)$$

$$n_s^-(0) = \frac{N_{2D}}{2} F_0(\eta_F - qV_D)$$

$$J^+ = \frac{qN_{2D}}{2} v_T [F_{1/2}(\eta_{F1})]$$

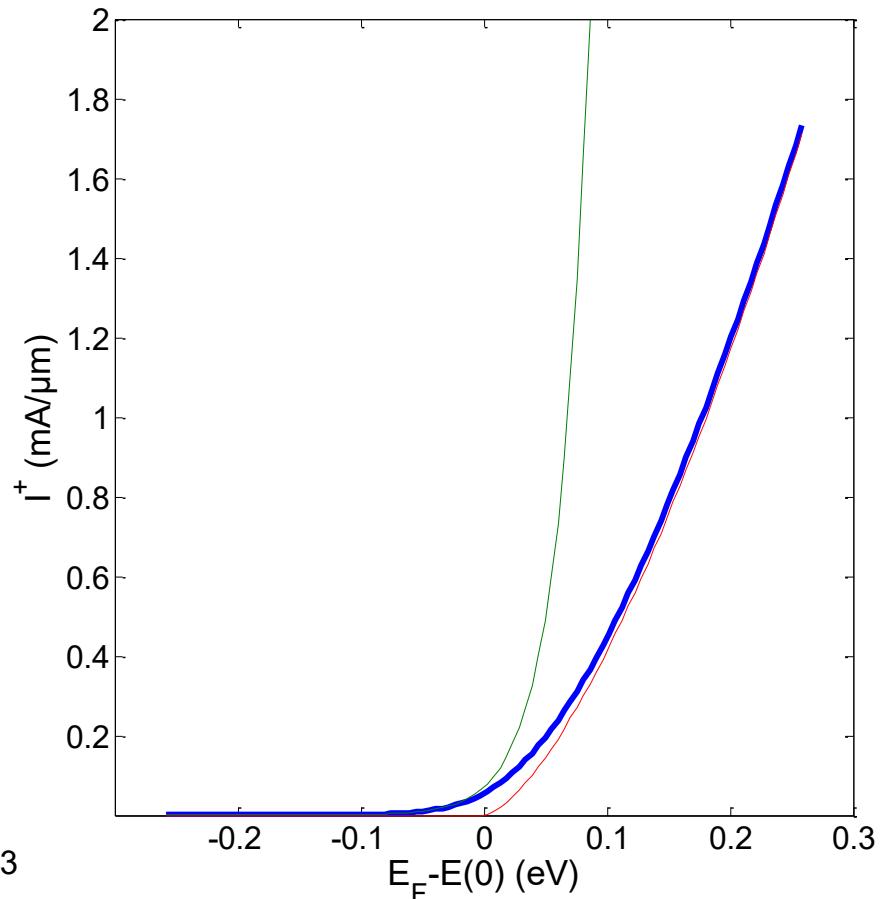
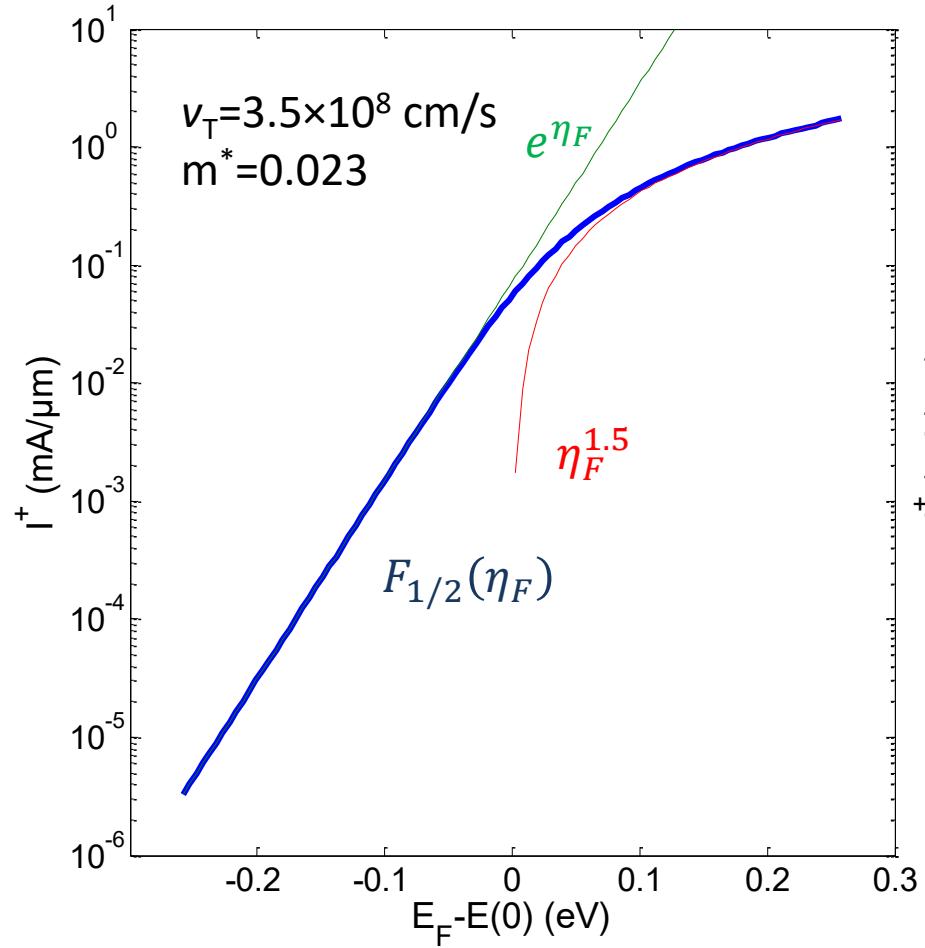
$$J^- = \frac{qN_{2D}}{2} v_T [F_{1/2}(\eta_{F1} - U_D)]$$

$$J^+ = q n_s^+ v^+$$

$$v^+ = v_T \frac{F_{1/2}(\eta_F)}{F_0(\eta_F)}; v_T = \sqrt{\frac{2kT}{\pi m^*}}$$

$$v^- = v_T \frac{F_{1/2}(\eta_F - U_D)}{F_0(\eta_F - U_D)}$$

I⁺: T=300K



Quantum Capacitance / Bias self consistency

$$V'_G = \psi_s + Q/C_{ox}$$

$$V'_G = -\frac{\varepsilon_1(0)}{q} + \frac{qn_s(\varepsilon_1)}{C_{ox}}$$

$$n_s = \frac{N_{2D}}{2} \{F_0(\eta_F) + F_0(\eta_F - U_D)\}$$

$$J^+ = \frac{qN_{2D}}{2} v_T [F_{1/2}(\eta_F)]$$

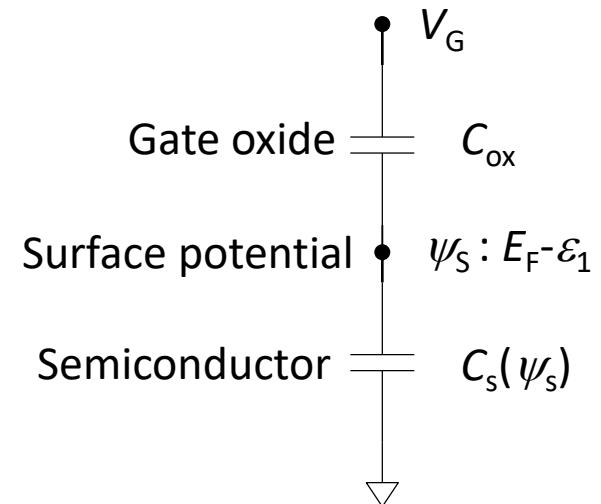
$$J^- = \frac{qN_{2D}}{2} v_T [F_{1/2}(\eta_F - U_D)]$$

MOS-limit: $C_{ox} \ll C_s$: n_s constant

Bipolar limit: $C_{ox} \rightarrow \infty$ $\delta\varepsilon_1 = -q\delta V_G$

Solve for η_F

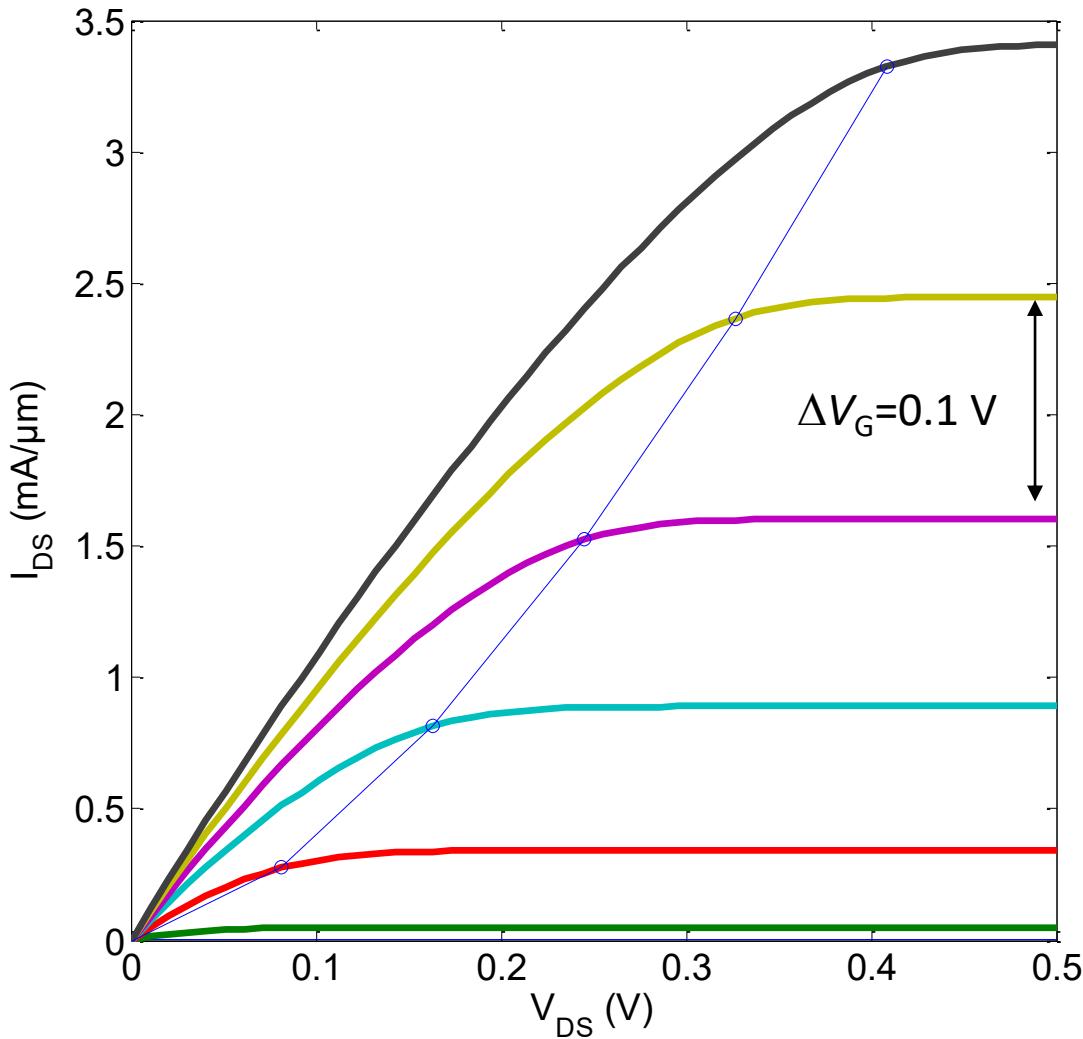
Calculate $I^+ - I^-$



$$\delta\varepsilon_1 = -\frac{q\delta V_G}{1 + \frac{C_s}{C_{ox}}}$$

If we also include the charge centroid $C_{ox} \rightarrow C_{ox} || C_C$ capacitance:

Modeled 2D I-V



$m^*=0.023/m_0$
 $t_{ox}=5 \text{ nm}$
 $\varepsilon_r=20$

Degenerate expressions with Quantum Capacitance

$$I^+ \approx \frac{qW2\sqrt{2m^*}}{3\pi^2\hbar^2} \left(\frac{qC_{ox}}{C_{ox} + C_s} \right)^{\frac{3}{2}} (V_{GS} - V_T)^{\frac{3}{2}}$$
$$C_s = \frac{C_q}{2} = \frac{m^*}{2\pi\hbar^2}$$

$$V_{ds,sat} \approx \frac{V_{GS} - V_T}{1 + \frac{C_q}{2C_{ox}}}$$

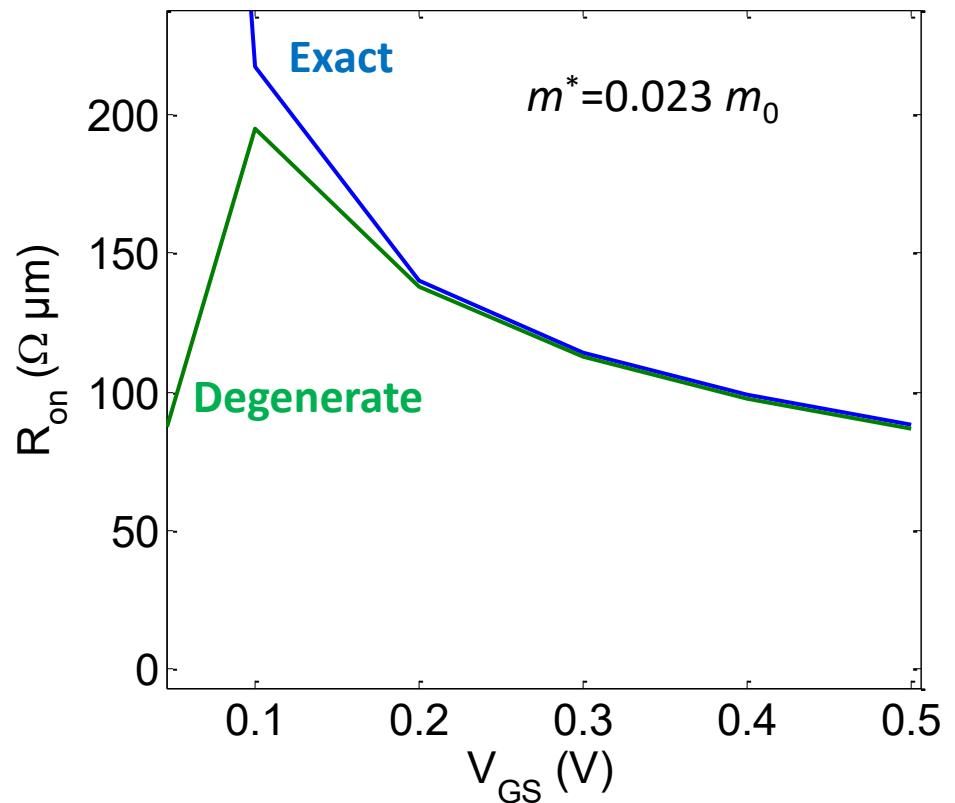
$$G_{CH} = \frac{q^2 W \sqrt{2m^*}}{\pi^2 \hbar^2} \sqrt{q \frac{C_{ox}}{C_{ox} + C_q} (V_G - V_T)}$$

Minimum on-resistance: $R_{on} = \frac{1}{G_{CH}}$

R_{on} - resistance

A ballistic FET/conductor has a minimum on-resistance!

This is important when trying to extract R_c from measured data.



Density of states bottle neck

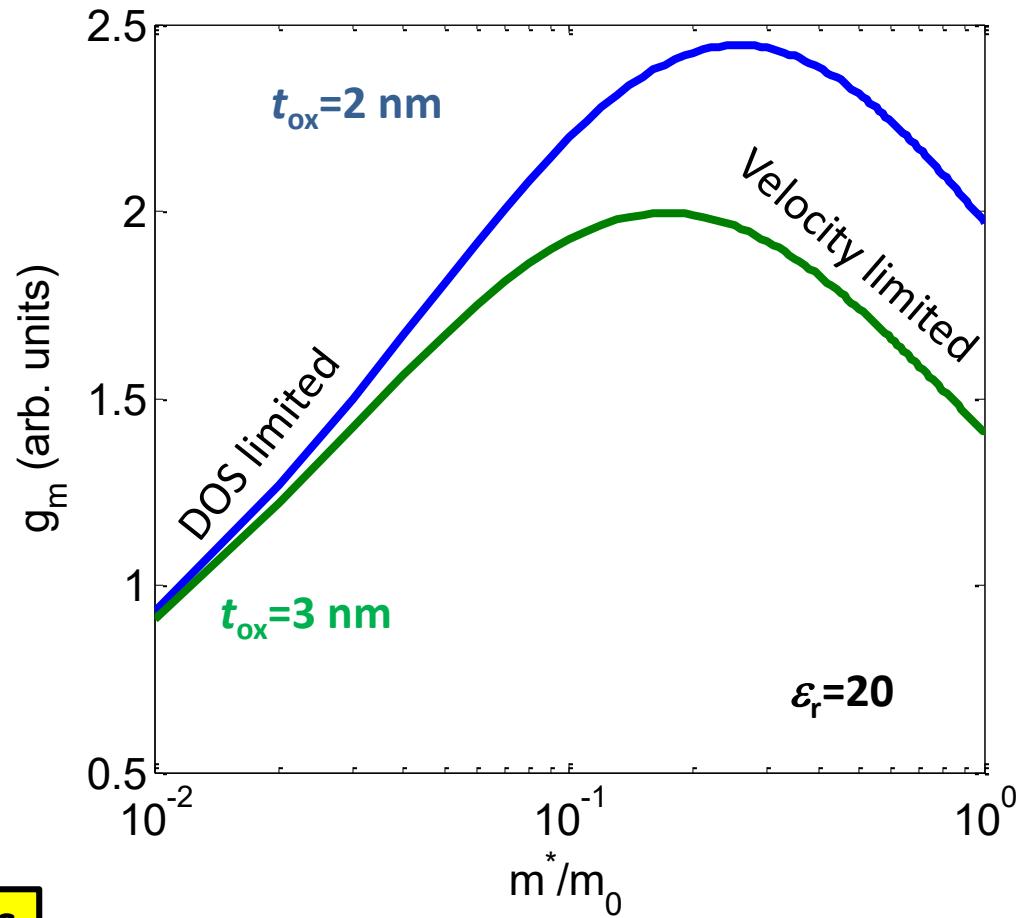
Optimal mass when:

$$\frac{C_q}{4} = C_{ox}$$

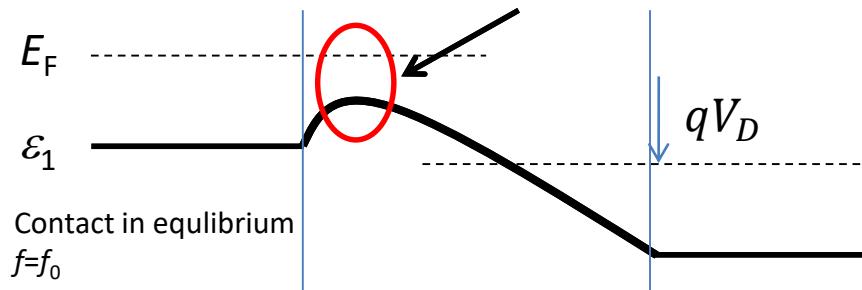
$$m_{opt}^* = \frac{4C_{ox}\pi\hbar^2}{q^2}$$

$$I^+ \propto \sqrt{m^*} \left(\frac{1}{1 + \frac{m^*}{2\pi\hbar^2 C_{ox}}} \right)^{\frac{3}{2}}$$

Scattering is proportional to DOS
Low m^* → less scattering



MOS statistics – MOS Limit



$$C_{ox} \ll C_S, C_q$$

$$n_s = n_s^+(E_F) + n_s^-(E_F, V_D) = C_{ox}(V_{GS} - V_T)/q$$

Allows us easily to solve for η_F
And to relate n_s to V_{GS}

Charge at the top of the barrier is constant,
independent of V_D

This is correct if $C_{ox} \ll C_s$ and above V_T

$E_F - \epsilon_1$ needs to increase as V_D increases

- }
1) No 2D effects
2) Large DOS, thick C_{ox}

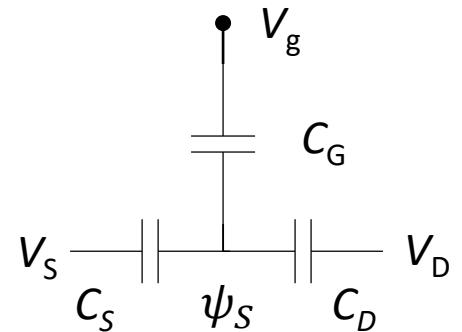
2D Electrostatics – ballistic FETs

$$V'_G = -\frac{\varepsilon_1(0)}{q} + \frac{qn_s(\varepsilon_1(0))}{C_{ox}} \quad 1D$$

$$\varepsilon_1(0) = U_L + U_P$$

$$U_L = \alpha_G V_G + \alpha_D V_D + \alpha_S V_S \quad \text{Laplace Eq.}$$

$$U_P = \frac{q^2}{C_\Sigma} n_s(\varepsilon_1(0)) = U_C n_s \quad \text{Mobile Charge/Possion Eq.}$$

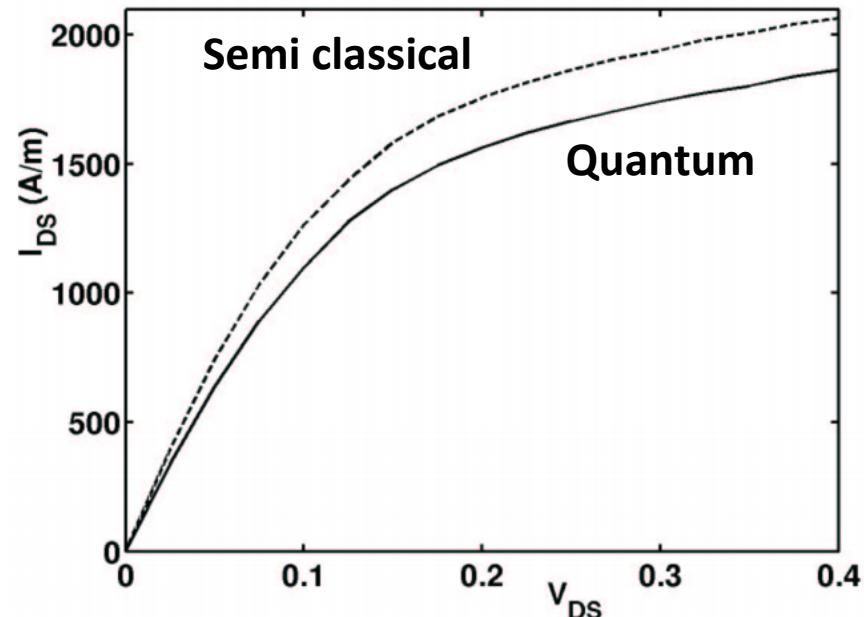
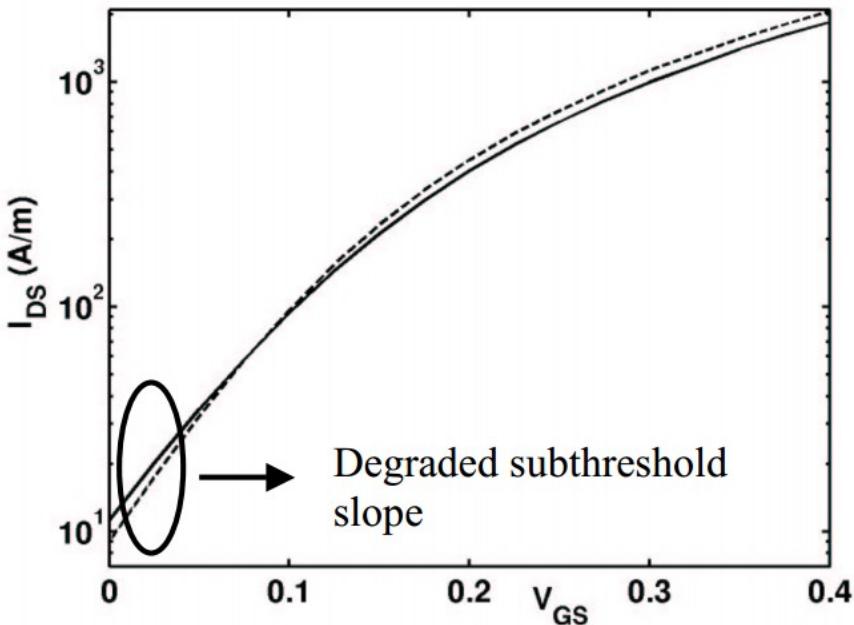


α_G : subthreshold slope, g_m

α_D : DIBL/ g_d

E_F : correct off-current

Quantum / Semiclassical



$$I^+ = \frac{2q}{h} \int_0^\infty M(E)T(E)f_0(E, E_F)dE$$

Semi-classical: We assume $T(E)$ step function

Quantum: $T(E) < 1$ above $E(0)$ due to reflections
 $T(E) > 0$ below $E(0)$ due to tunneling

Quantum effects are small for $L_g=10$ nm (Si) devices!