



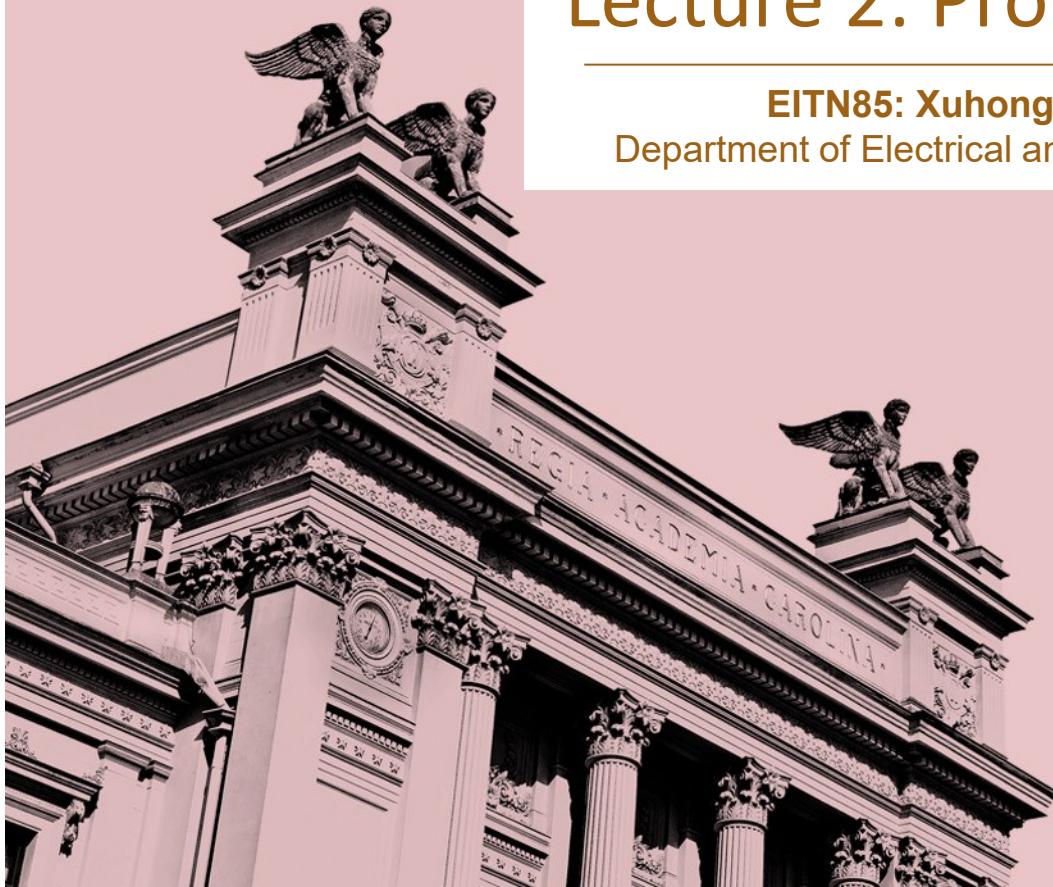
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Wireless Communications Channels

Lecture 2: Propagation Mechanisms

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Recap: Last Lecture

Contents

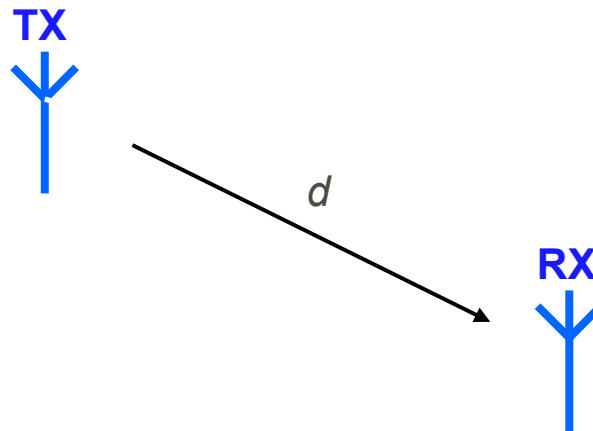
Propagation mechanisms:

- ☐ Free space attenuation
- ☐ Reflection and transmission
- ☐ Diffraction
- ☐ Diffuse scattering
- ☐ Waveguiding

Examples from real world propagation scenarios

Free space attenuation

- ❑ Assume TX and RX antennas in **free space** and would like to derive the received power as a function of link distance and transmit power
 - ❑ Assume omnidirectional antennas for now

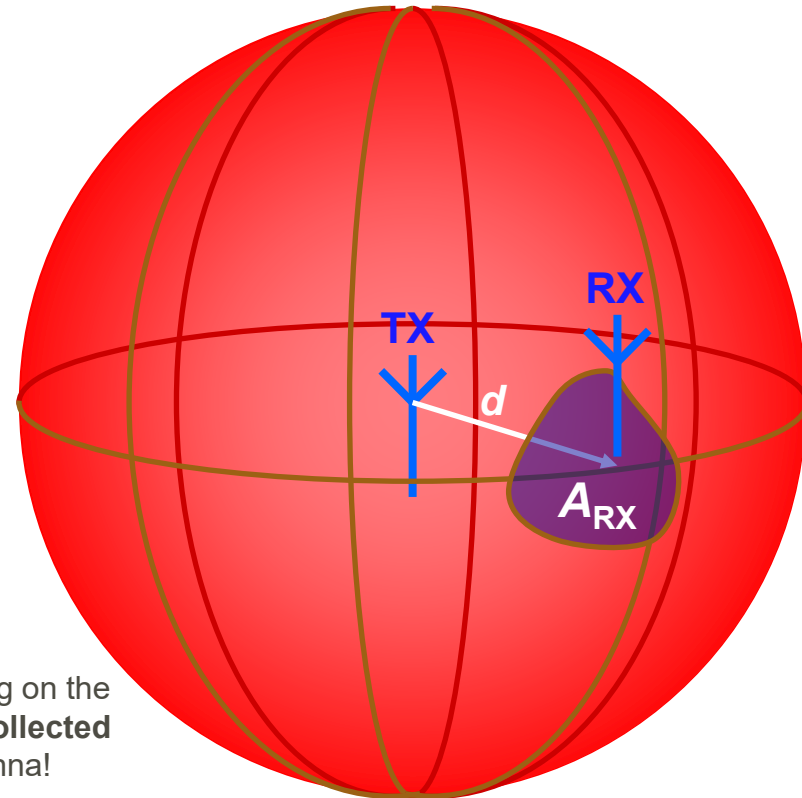


Free space attenuation

- ❑ Assume TX and RX antennas in **free space** and would like to derive the received power as a function of link distance and transmit power (omnidirectional antennas)
- ❑ Energy conservation: integral of power density over any **closed surface** = transmit power
- ❑ If TX antenna radiates isotropically, then power density on surface is $P_{TX}/(4\pi d^2)$. Then,

$$P_{RX}(d) = P_{TX} \frac{1}{4\pi d^2} A_{RX}$$

Power impinging on the area which is **collected** by the RX antenna!



Sphere of radius d



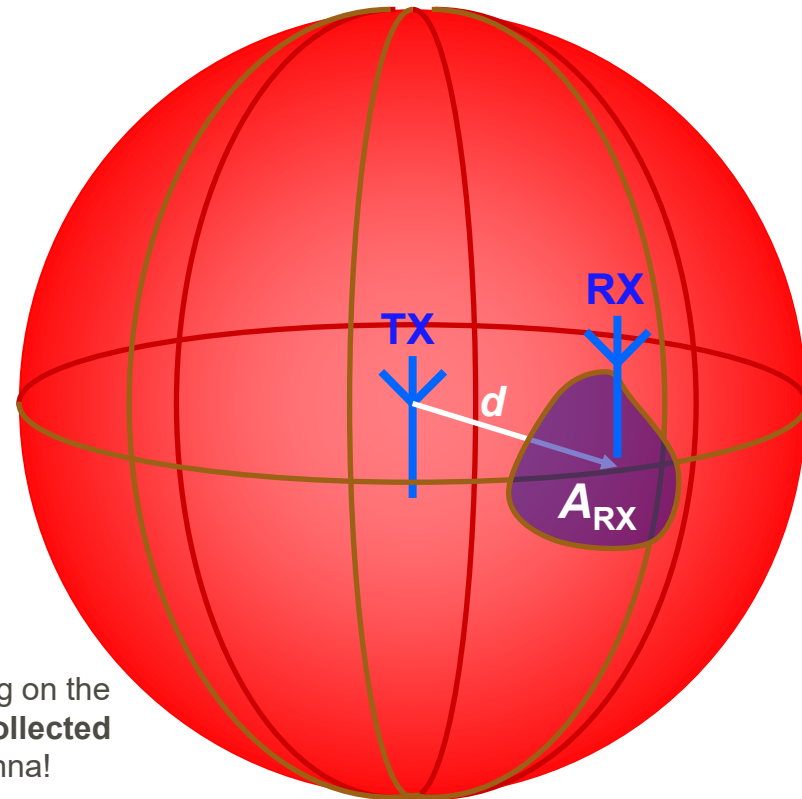
Free space attenuation

- ❑ Assume TX and RX antennas in **free space** and would like to derive the received power as a function of link distance and transmit power
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- ❑ Product of TX power and gain is known as: effective isotropic radiated power (EIRP)

- ❑ Relationship between effective area and antenna gain: $G_{RX} = \left(\frac{4\pi}{\lambda^2}\right) A_{RX}$



Shpere of radius d



Free space loss: Friis' law

Received power, with antenna gains G_{TX} and G_{RX} :

$$P_{RX}(d) = \frac{G_{RX} G_{TX}}{L_{free}(d)} P_{TX} = P_{TX} \left(\frac{\lambda}{4\pi d} \right)^2 G_{RX} G_{TX} *$$

Free space loss factor

RX power goes down as a function of frequency, for a fixed distance.

$$\begin{aligned} P_{RX|dB}(d) &= P_{TX|dB} + G_{TX|dB} - L_{free|dB}(d) + G_{RX|dB} \\ &= P_{TX|dB} + G_{TX|dB} - 10 \log_{10} \left(\frac{4\pi d}{\lambda} \right)^2 + G_{RX|dB} \end{aligned}$$

Question: What happens if d is 0 in $*$?



Free space loss: Friis' law implications

$$P_{RX}(d) = \frac{G_{RX} G_{TX}}{L_{free}(d)} P_{TX} = \boxed{P_{TX} \left(\frac{\lambda}{4\pi d} \right)^2 G_{RX} G_{TX}}$$

Free space loss:

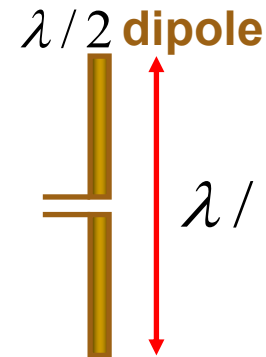
What and where is the far field?

The free space loss calculations are only valid in the "far field" of the antennas.

Far-field conditions are assumed far beyond the "Rayleigh" distance:

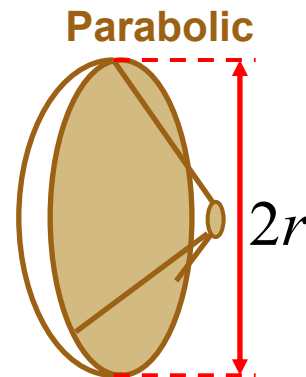
$$d_R = \frac{2L_a^2}{\lambda}$$

where L_a is the largest dimension of the antenna.



$$L_a = \lambda/2$$

$$d_R = \lambda/2$$



$$L_a = 2r$$

$$d_R = \frac{8r^2}{\lambda}$$

Another rule of thumb is: "At least 10 wavelengths"

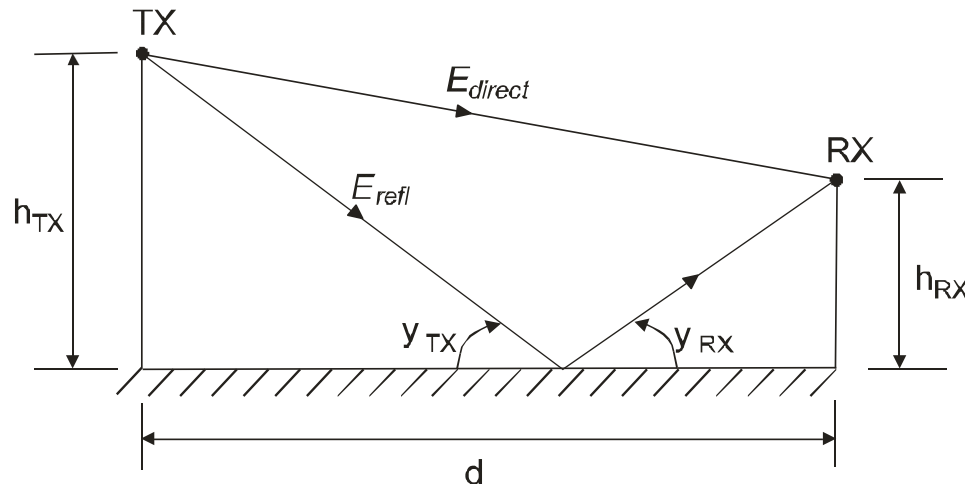
Quiz

Compute the Rayleigh distance of a square patch antenna receiving a signal with a gain of 10 dB.



The d^{-4} law – I

Instead of just considering a direct path, let's look at the following scenario



the power behaves as

$$P_{RX}(d) \approx P_{TX} G_{TX} G_{RX} \left(\frac{h_{TX} h_{RX}}{d^2} \right)^2$$

for distances greater than

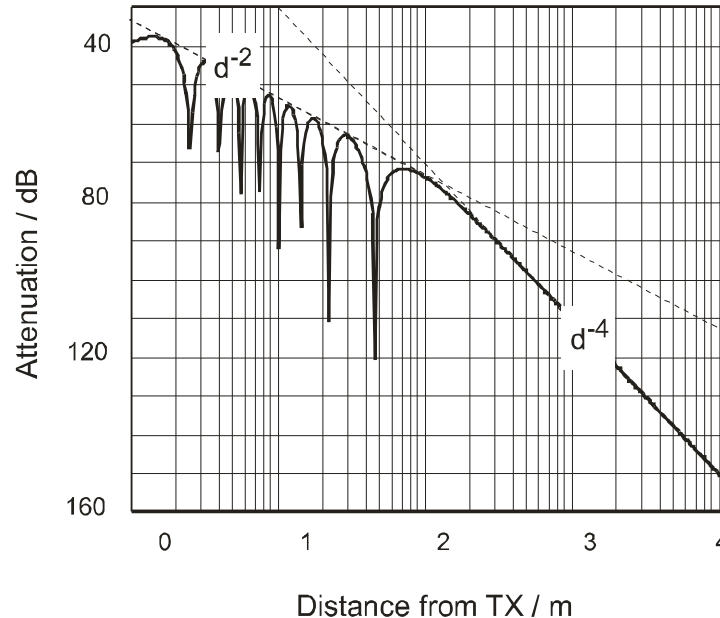
$$d_{break} \gtrsim 4h_{TX}h_{RX}/\lambda$$



Continuation of Slide (10): How?

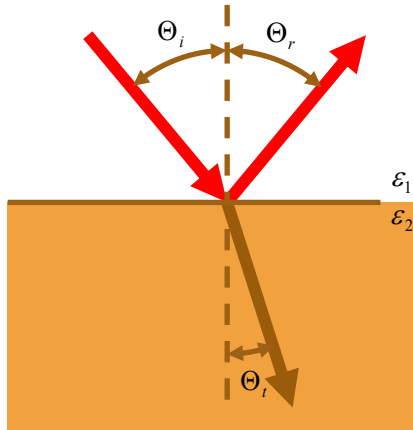


The d^{-4} law – II

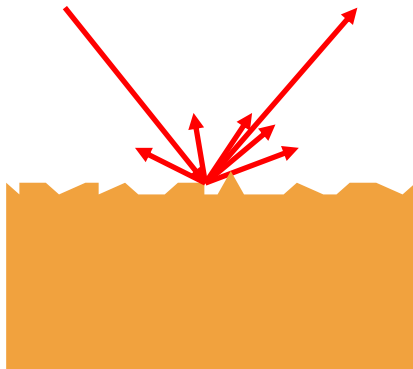


- However
 - $n=4$ is not a universal decay exponent
 - Theoretical model is not fulfilled in practice
 - Breakpoint is rarely where theoretically predicted
 - Second breakpoint at the radio horizon

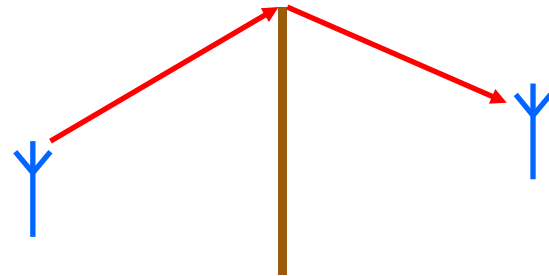
Fundamental propagation mechanisms



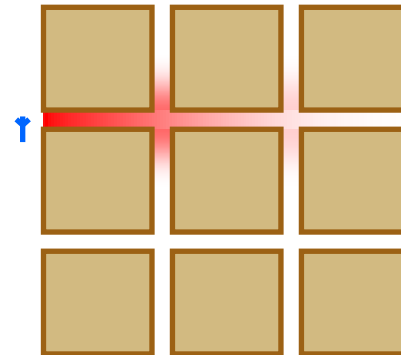
Reflection and transmission



Scattering



Diffraction



Waveguiding

Complex dielectric constant

Lets take a homogeneous planewave incident onto a dielectric half-space

$$\delta_i = \epsilon_i - j \frac{\sigma_{e,i}}{2\pi f_c}$$

dielectric constant, permittivity

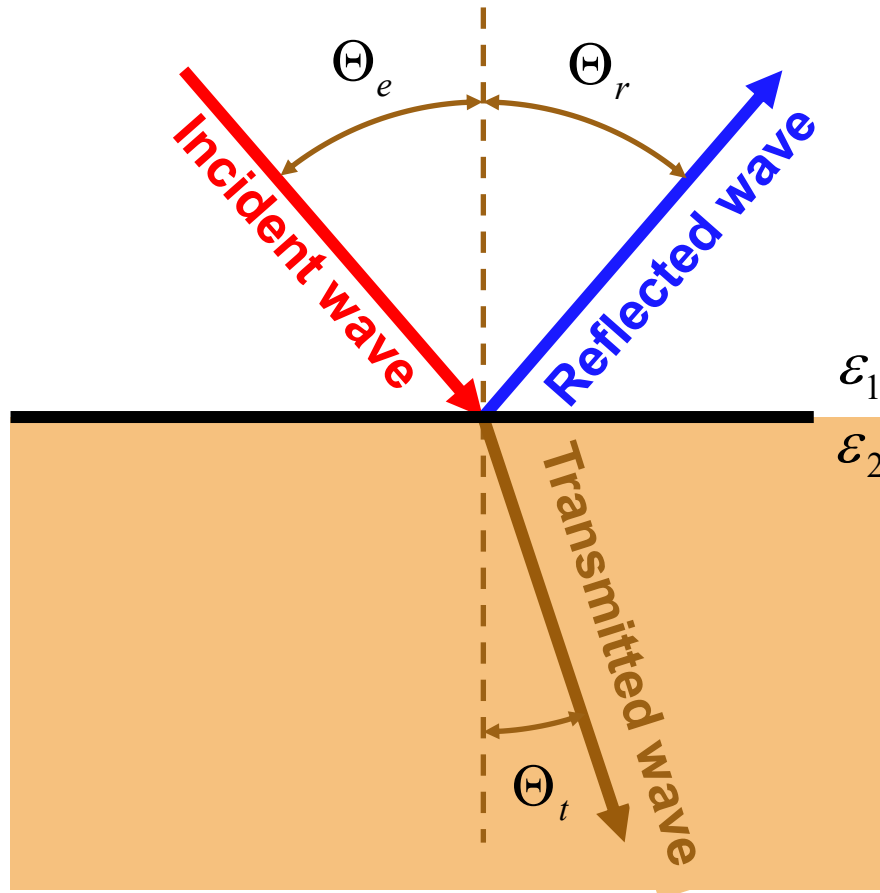
conductivity

Describes the dielectric material in one single parameter

Examples	Permittivity	conductivity
Concrete	6	10^{-2}
Gypsum	6.5	10^{-2}
Wood	23	10^{-11}
Glass	5	10^{-12}
Air	1	



Reflection and transmission



Reflected angle:

$$\Theta_e = \Theta_r.$$

Transmitted angle:

$$\frac{\sin \Theta_t}{\sin \Theta_e} = \frac{\sqrt{\epsilon_1}}{\sqrt{\epsilon_2}}$$



What is Reflection and Transmission Dependent on?

TM and TE Component Behaviors

Reflection
coefficient

$$\rho_{TM} = -\frac{\sqrt{\delta_2} \cos \Theta_e - \sqrt{\delta_1} \cos(\Theta_t)}{\sqrt{\delta_2} \cos \Theta_e + \sqrt{\delta_1} \cos(\Theta_t)}$$

$$\rho_{TE} = \frac{\sqrt{\delta_1} \cos(\Theta_e) - \sqrt{\delta_2} \cos(\Theta_t)}{\sqrt{\delta_1} \cos(\Theta_e) + \sqrt{\delta_2} \cos(\Theta_t)}$$

TM-waves



TE-waves



Transmission
coefficient

$$T = \sqrt{1 - \rho^2}$$



Transmission through layered structures

Total transmission coefficient

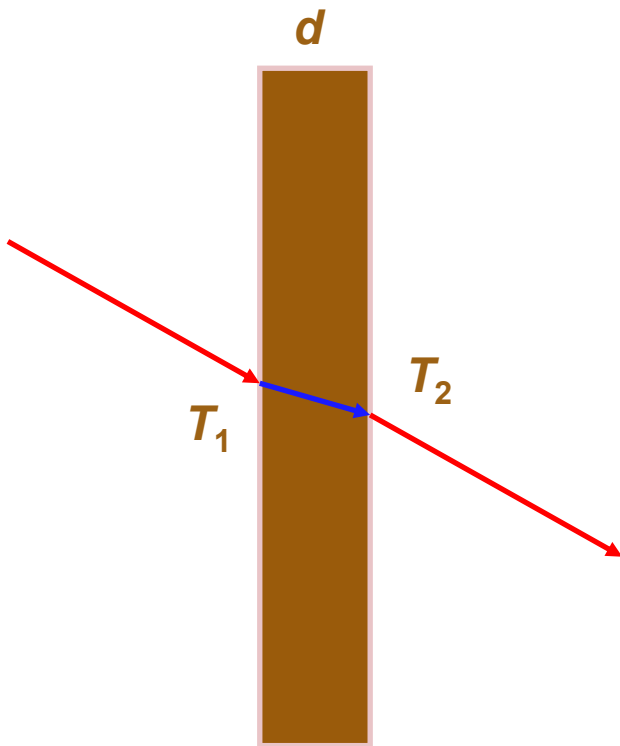
$$T = \frac{T_1 T_2 e^{-j\alpha}}{1 + R_1 R_2 e^{-2j\alpha}}$$

total reflection coefficient

$$\rho = \frac{\rho_1 + \rho_2 e^{-j2\alpha}}{1 + \rho_1 \rho_2 e^{-2j\alpha}}$$

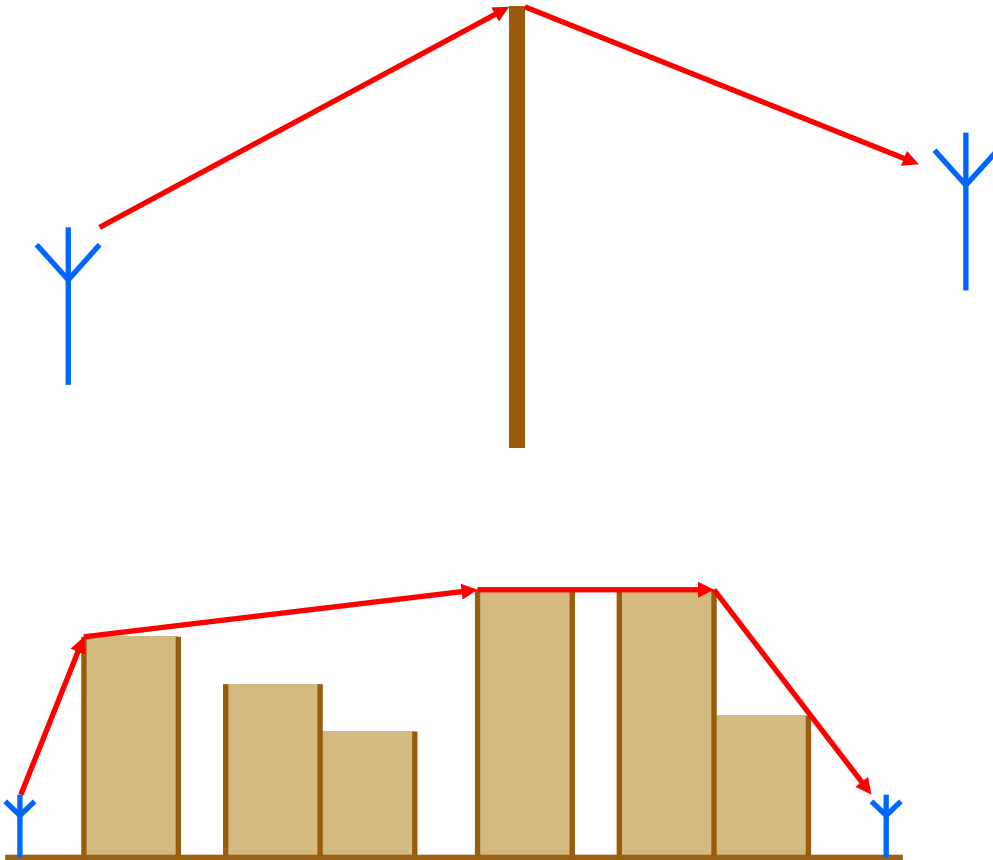
with the electrical length in the wall

$$\alpha = \frac{2\pi}{\lambda} \sqrt{\epsilon_1} d_{\text{layer}} \cos(\Theta_t)$$



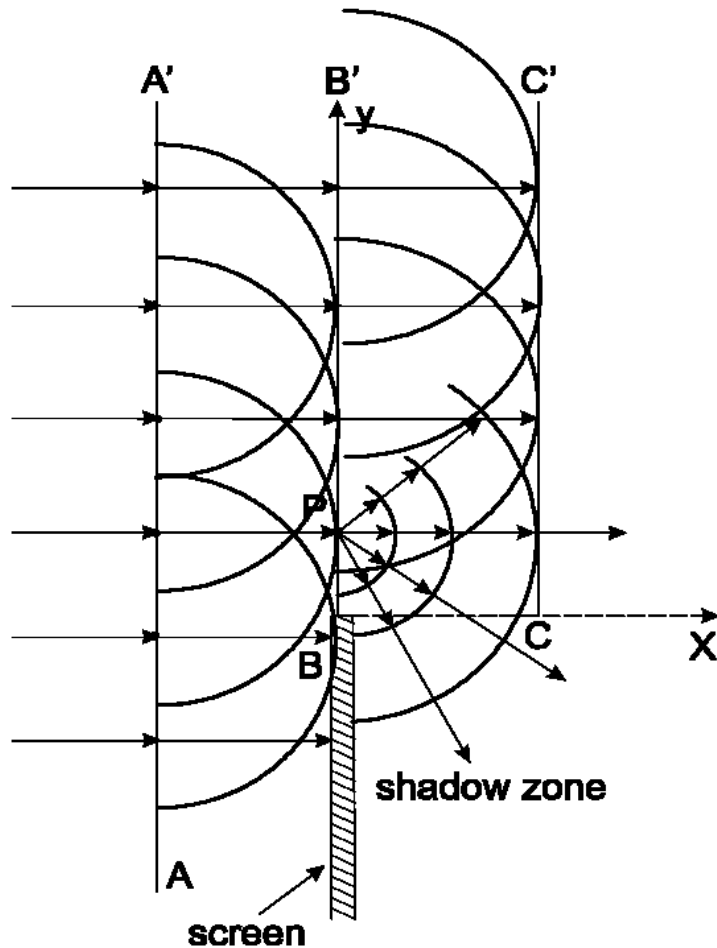
Wall with thickness d and two dielectrics

Diffraction: The principle



- Single or multiple edges
- makes it possible to go behind corners
- less pronounced when the wavelength is small compared to objects

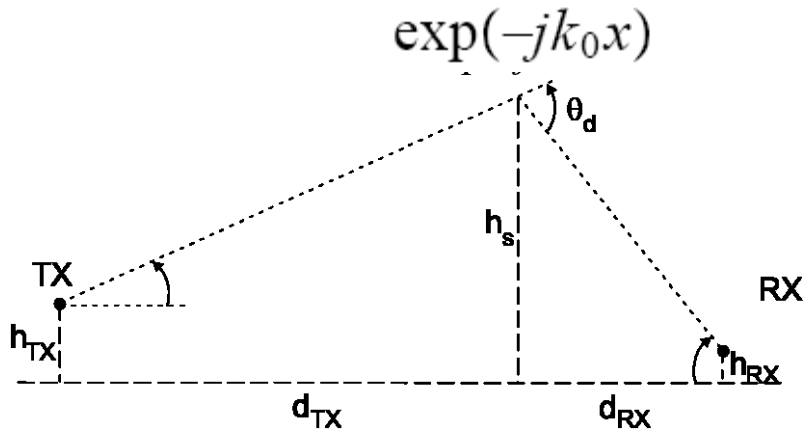
Diffraction: Huygen's principle



**Each point of a wavefront
can be considered as a
source of a spherical wave**

**➡ Bending around
corners and edges**

Diffraction coefficient



The Fresnel integral is defined

$$F(v_F) = \int_0^{v_F} \exp\left(-j\pi \frac{t^2}{2}\right) dt.$$

Total field

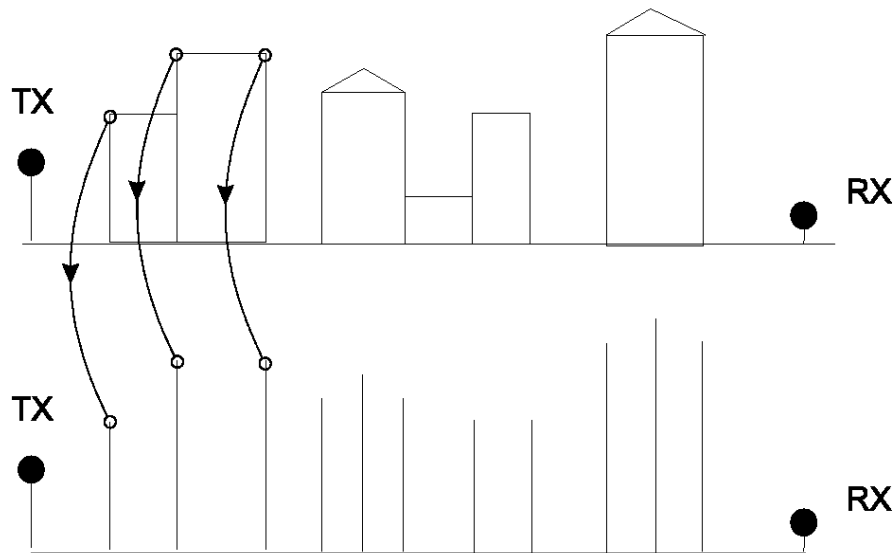
$$E_{\text{total}} = \exp(-jk_0 x) \left(\frac{1}{2} - \frac{\exp(-j\pi/4)}{\sqrt{2}} F(v_F) \right)$$

Fresnel integral

with the Fresnel parameter

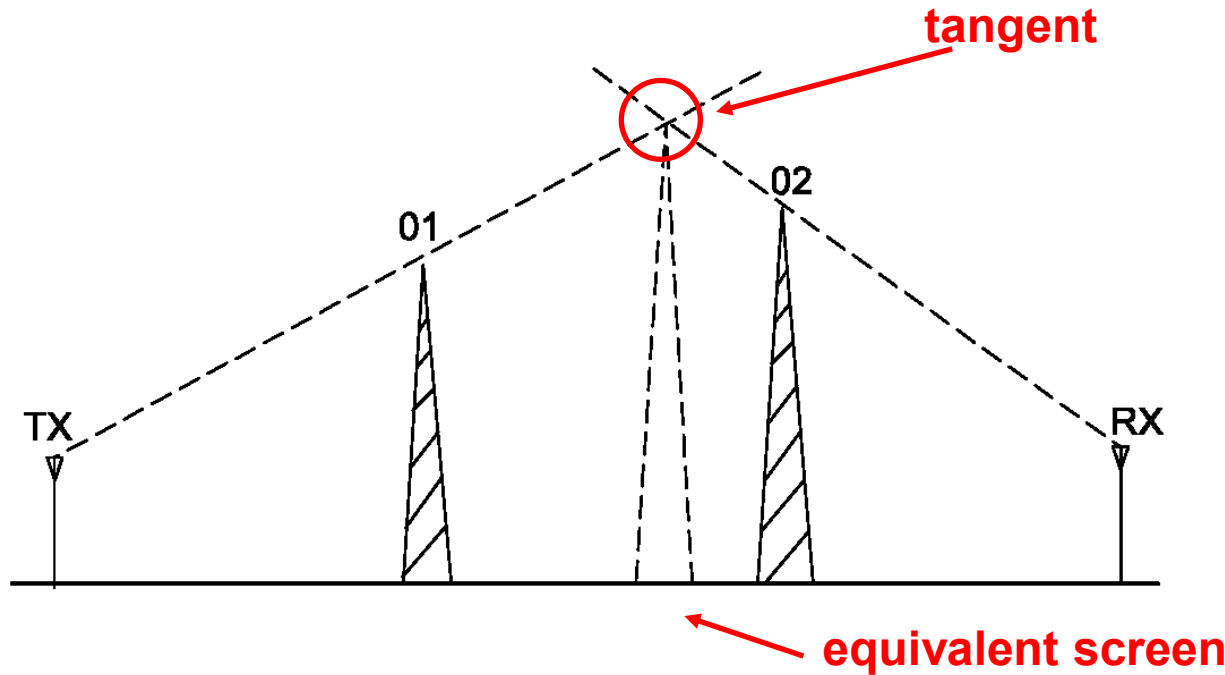
$$v_F = \alpha_k \sqrt{\frac{2d_1 d_2}{\lambda(d_1 + d_2)}}$$

Diffraction in real environments



For real environments we can represent buildings and objects as multiple screens

Diffraction: Bullington's method



Replace all screens with one equivalent screen

Height determined by the steepest angle

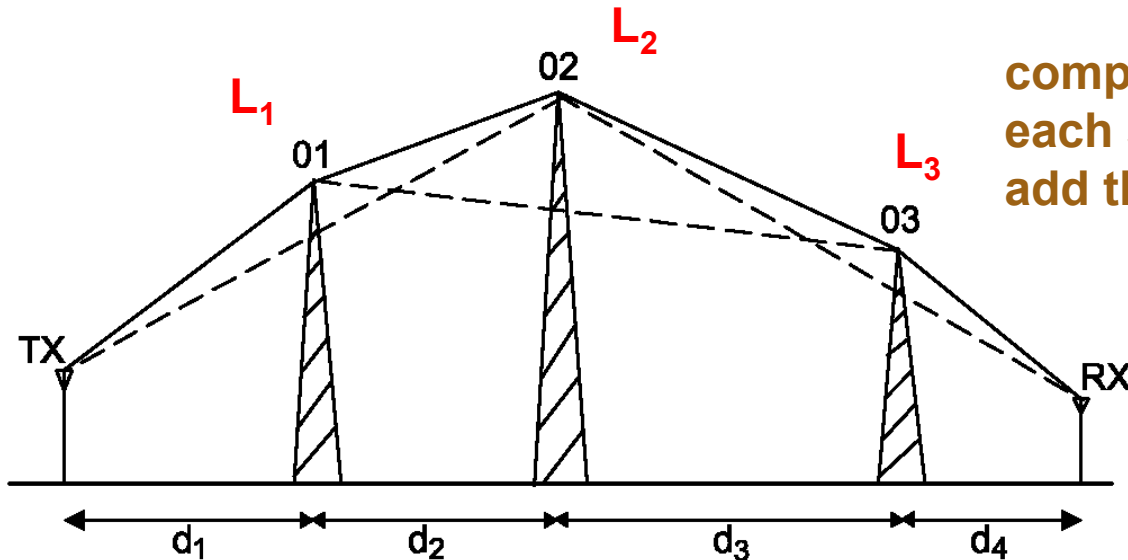
Simple but a bit optimistic

$$E_{\text{total}} = \exp(-jk_0x) \left(\frac{1}{2} - \frac{\exp(-j\pi/4)}{\sqrt{2}} F(v_F) \right)$$

$$v_F = \alpha_k \sqrt{\frac{2d_1d_2}{\lambda(d_1+d_2)}}$$



Diffraction – Epstein-Petersen Method



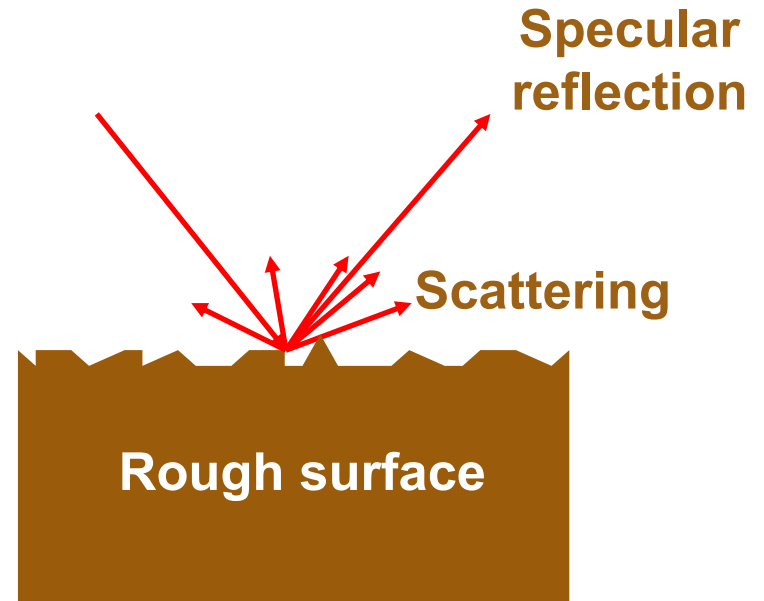
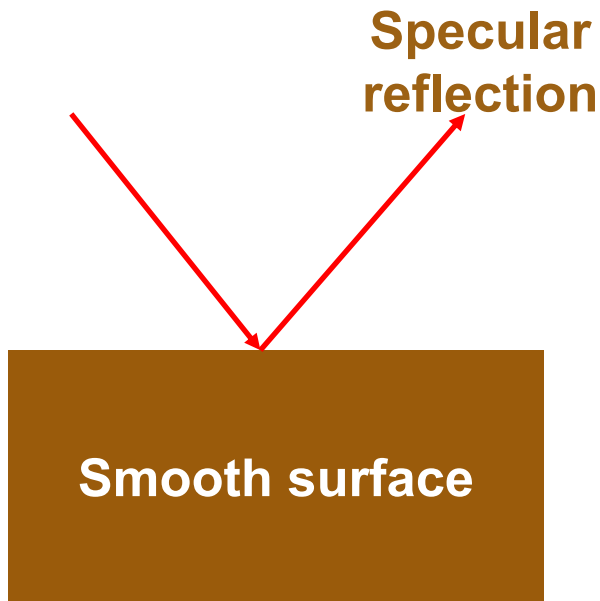
compute diffraction loss for each screen separately and add the losses

$$L_{\text{tot}} = L_1 + L_2 + L_3$$

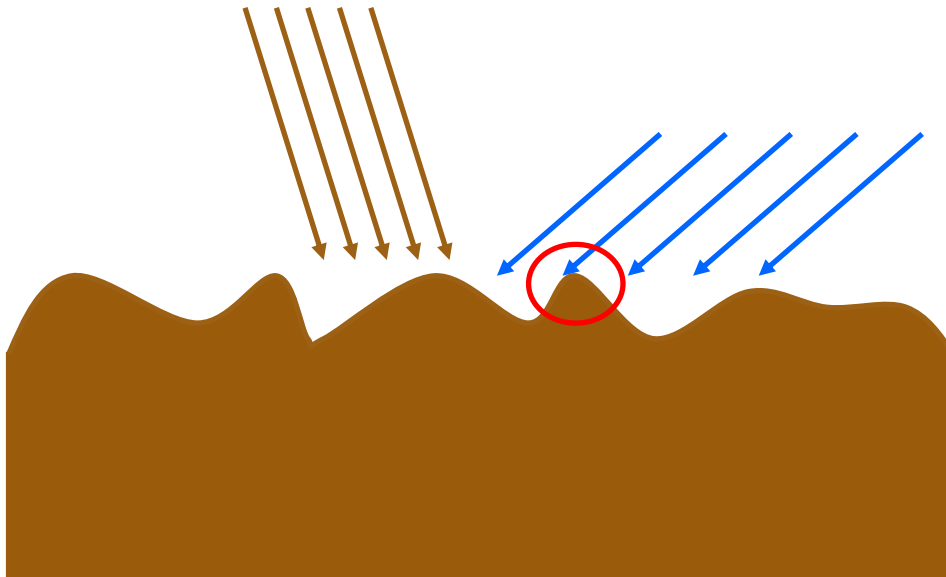
The same approach is used also for the ITU model, but with an empirical correction factor



Diffuse Scattering



Kirchhoff theory – scattering by rough surfaces



calculate distribution of the surface amplitude

assume no “shadowing” from surface

calculate a new reflection coefficient

for Gaussian surface distribution **angle of incidence**

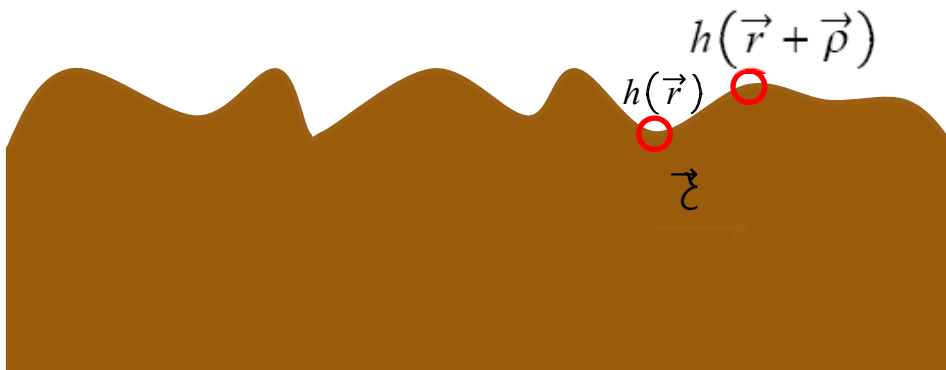
$$\rho_{\text{rough}} = \rho_{\text{smooth}} \exp \left[-2 \left(k_0 \sigma_h \sin \psi \right)^2 \right]$$

standard deviation of height



Perturbation theory – scattering by rough surfaces

$$\sigma_h^2 W(\vec{\rho}) = E_{\vec{r}} \{ h(\vec{r}) h(\vec{r} + \vec{\rho}) \}$$



Include shadowing effects by the surface

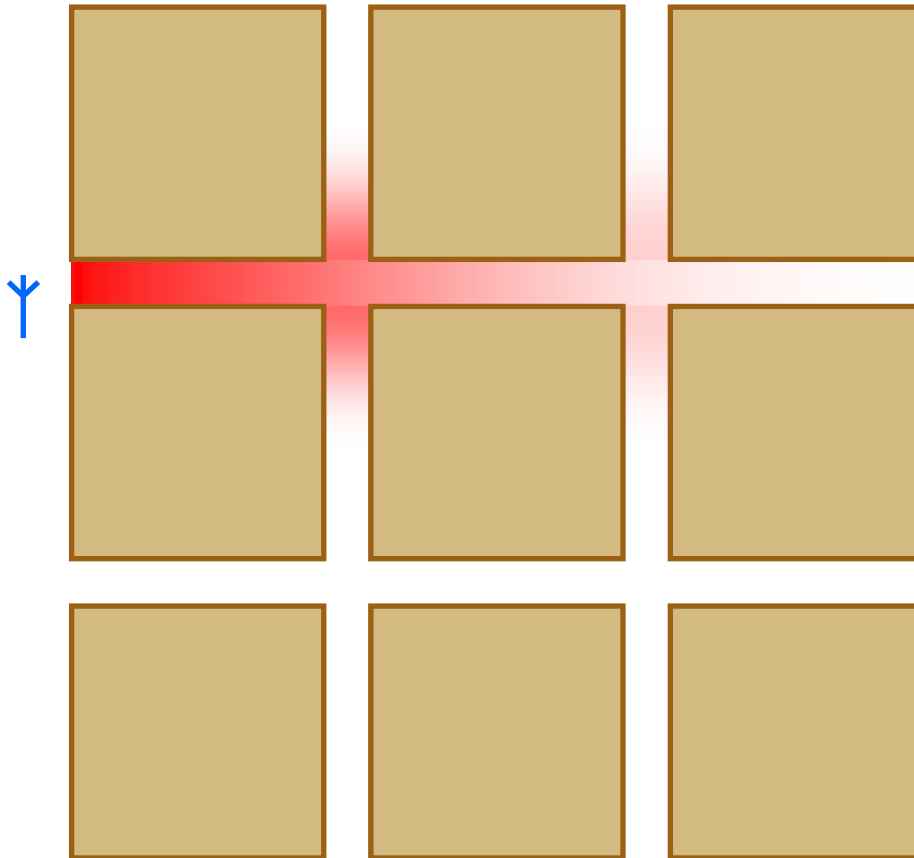
includes spatial correlation of surface – how fast are the changes in height

based on calculation of an “effective” dielectric constant

More accurate than Krichhoff theory, especially for large angles of incidence and “rougher” surfaces



Waveguiding



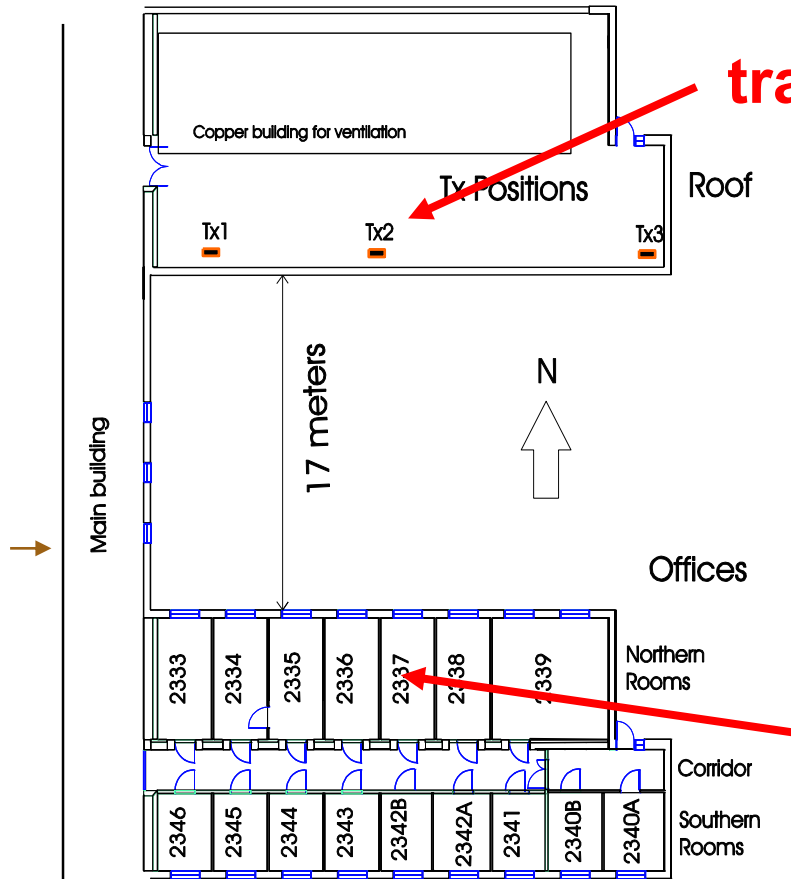
**Waveguiding effects
often result in lower
propagation
exponents**

$$n = 1.5-5$$

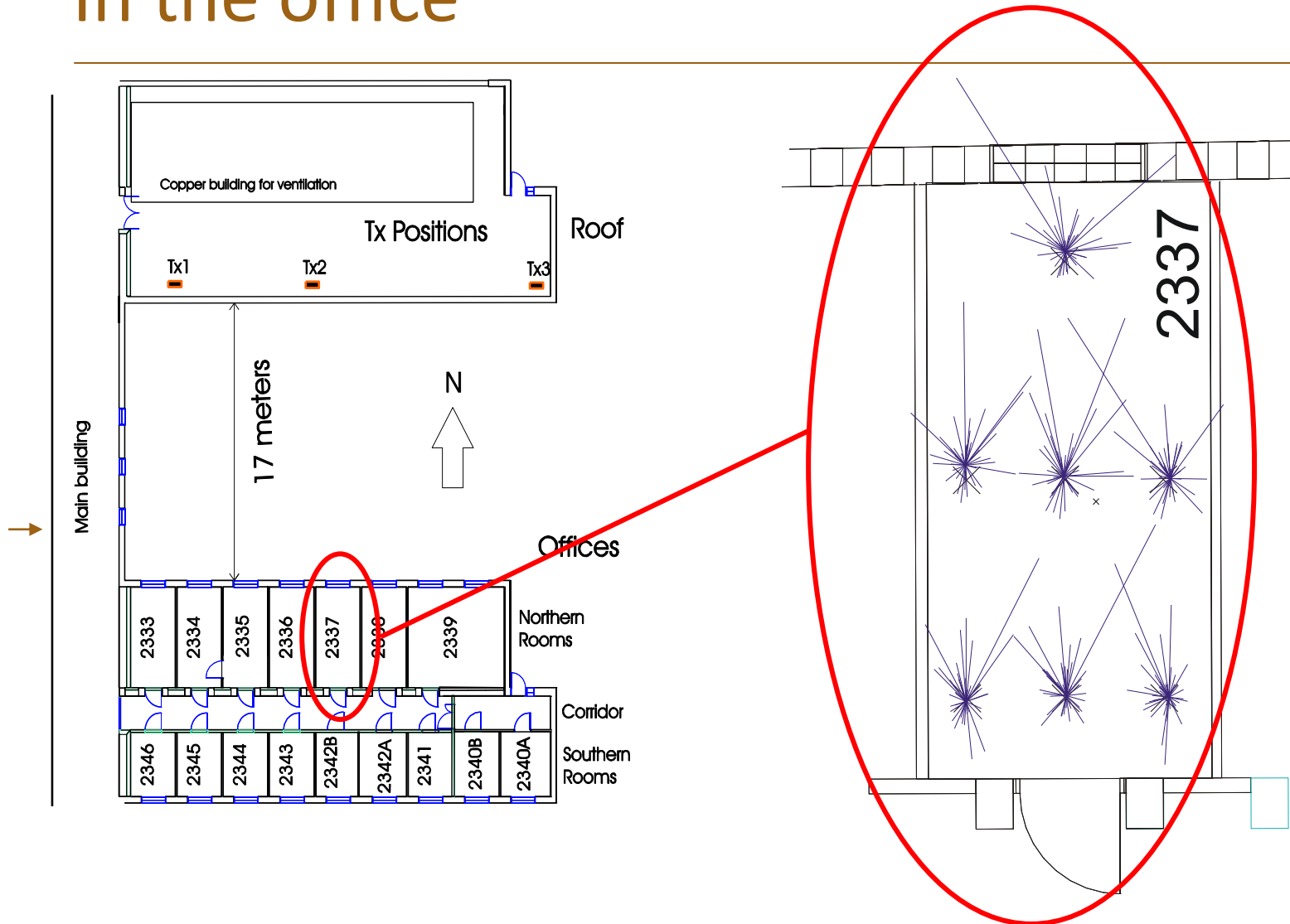
**This means lower
path loss along
certain street
corridors**



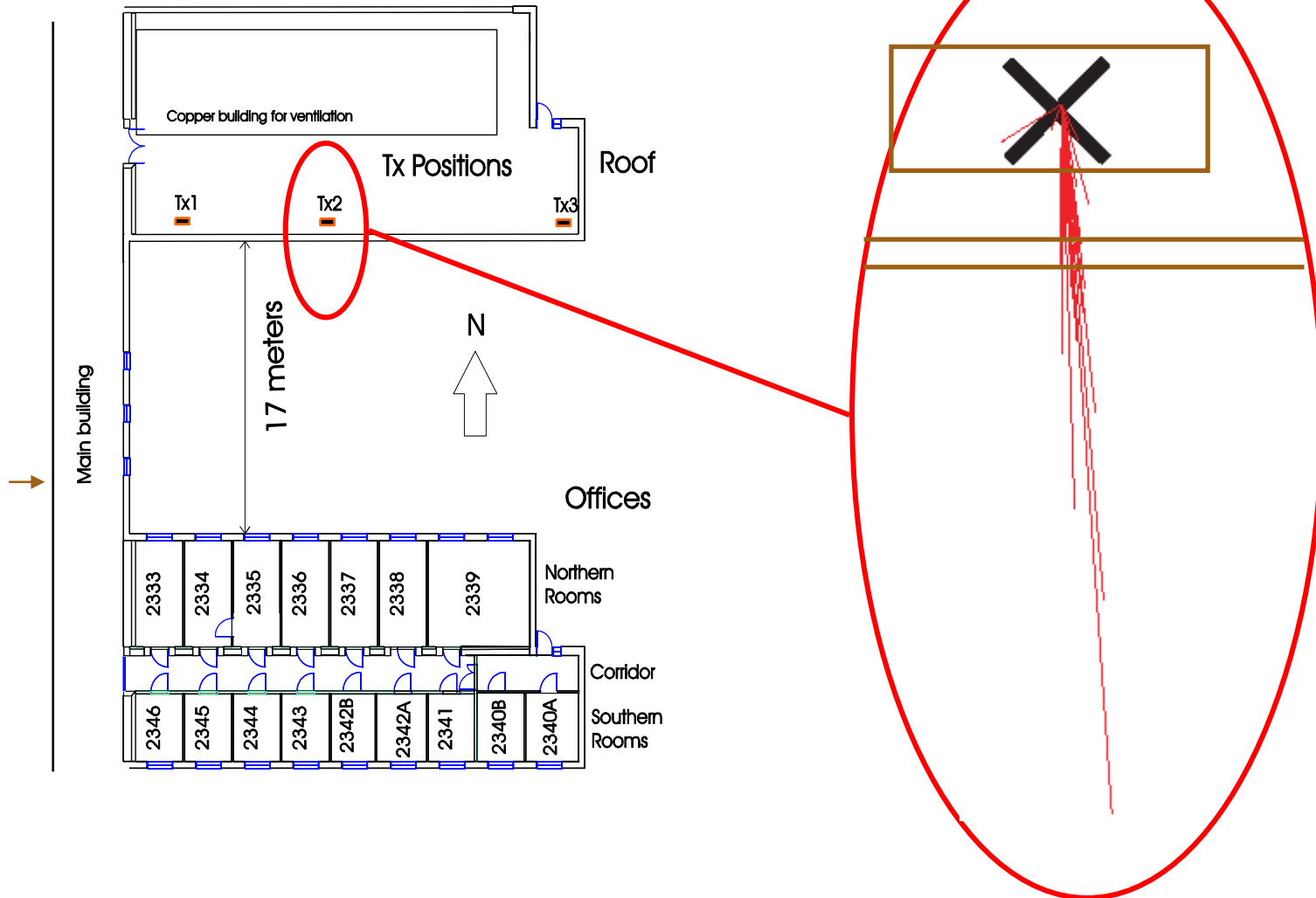
How does the signal reach the receiver Outdoor-to-indoor



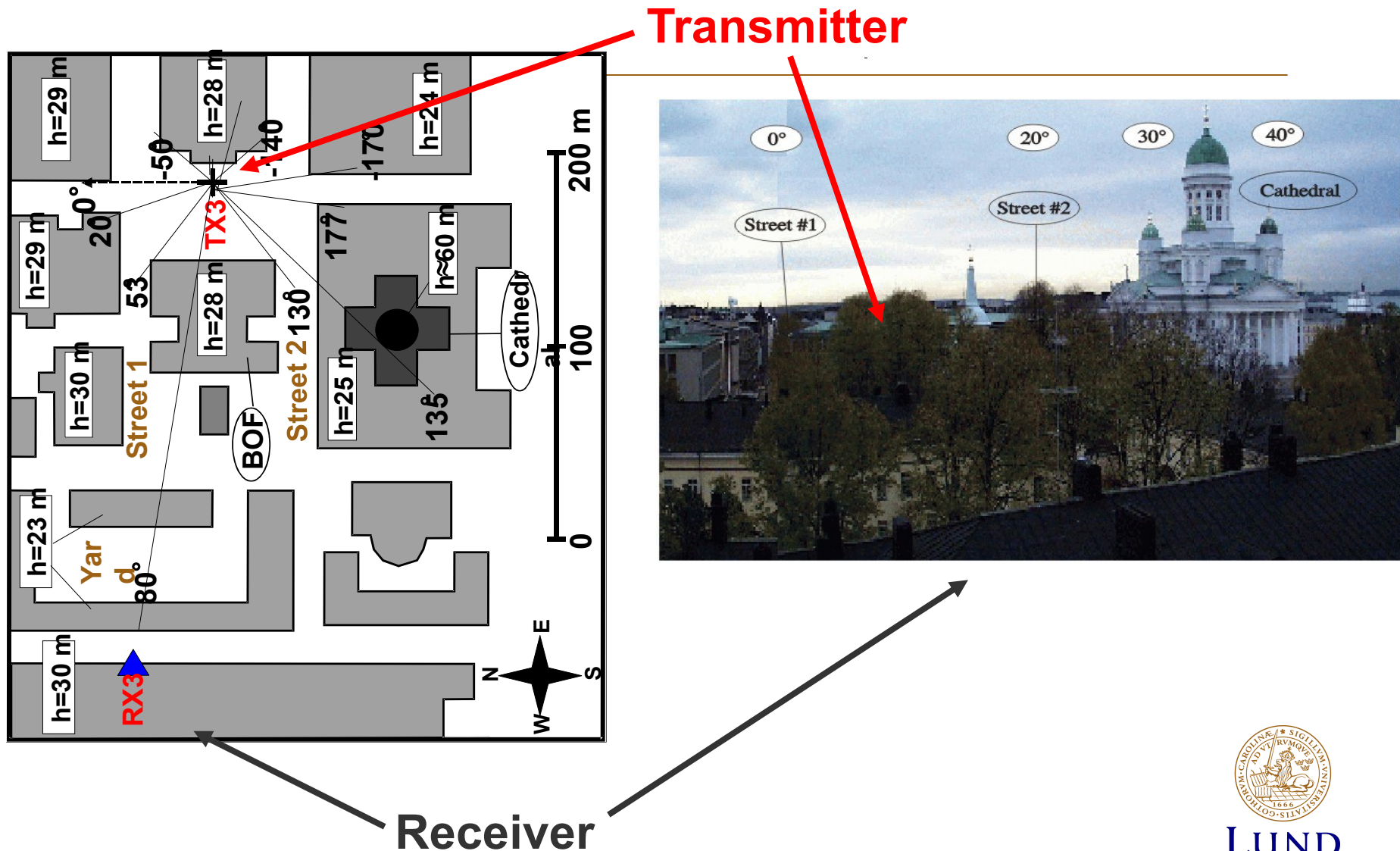
How does the signal reach the receiver In the office



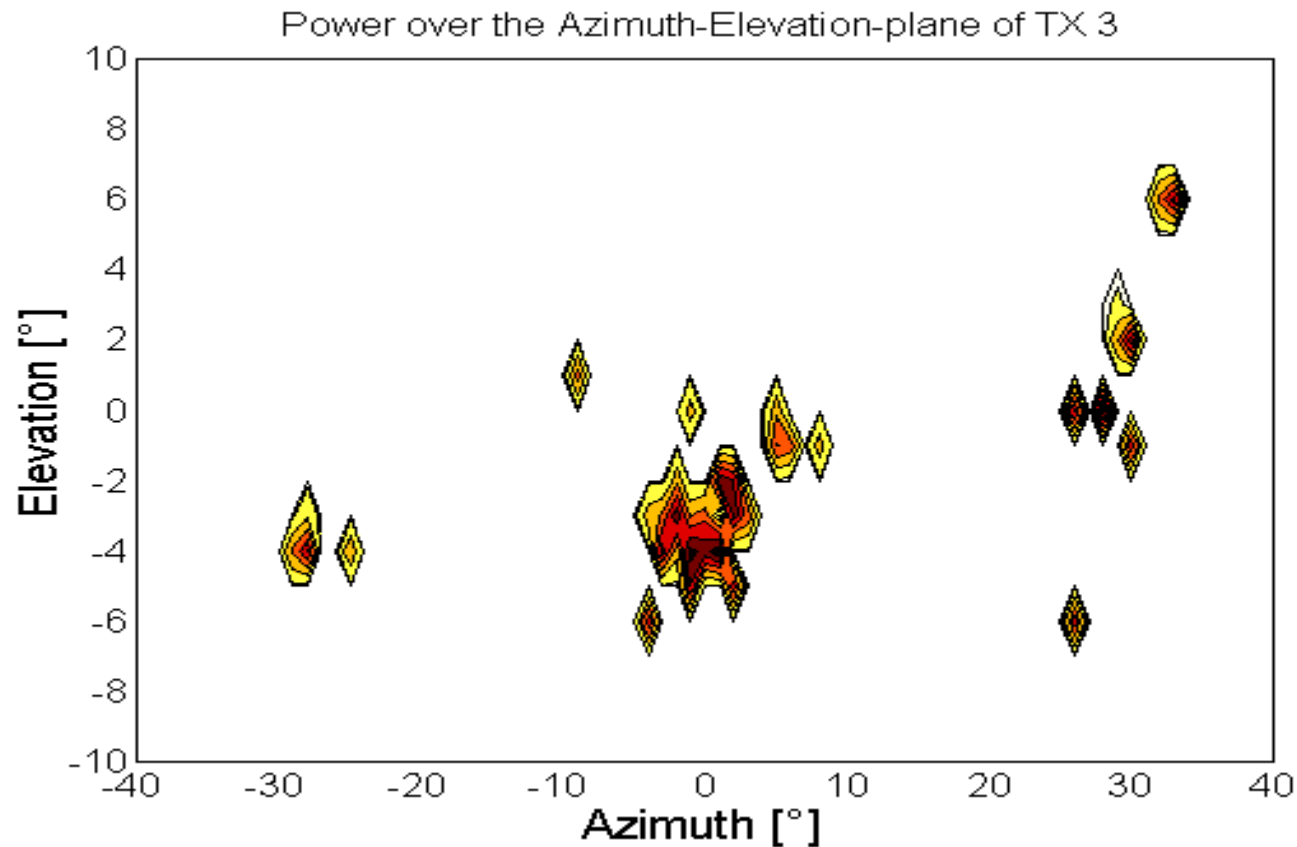
How does the signal leave the transmitter at the roof



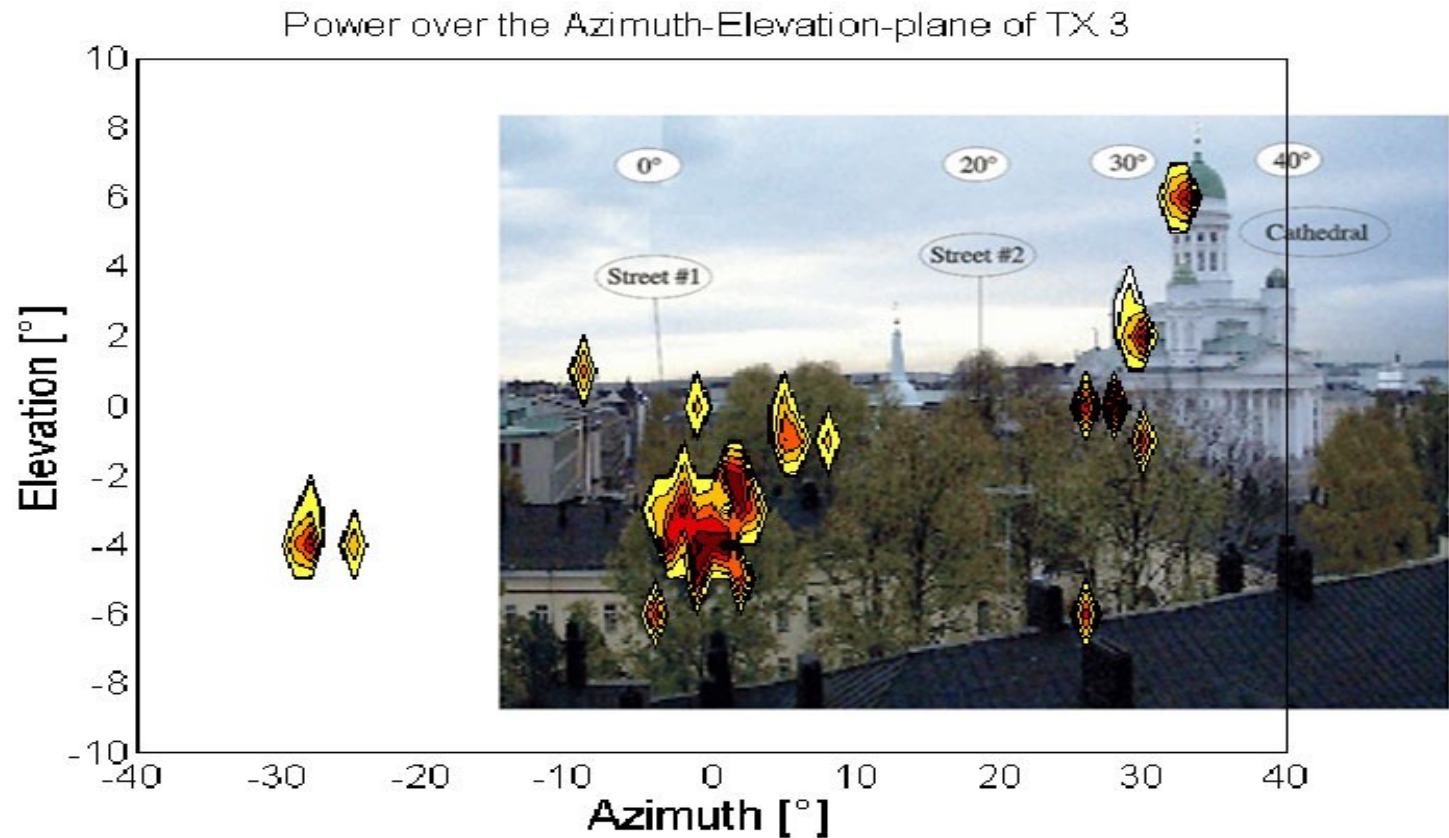
How does the signal reach the receiver outdoor urban



Signal arrives from some specific areas



Diffraction, reflection, scattering, transmission





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