



Electrodynamics, lecture 8, 2019

Anders Karlsson, anders.karlsson@eit.lth.se

Electrical and information technology

Visits to ESS and MAXIV

- ▶ ESS: May 9, 10:00 till 12:15
- ▶ MAX IV: May 23, 13.15-14.30

Last lecture

- ▶ Liénard-Wiechert potentials
- ▶ History of electromagnetism
- ▶ \mathbf{E} and \mathbf{B} from charged particle
- ▶ \mathbf{E} and \mathbf{B} from charged particle with constant velocity

Liénard-Wiechert potentials

$$\mathbf{A}(\mathbf{r}, t) = \frac{\mu_0}{4\pi} \frac{q\mathbf{v}(t_r)}{|\mathbf{r} - \mathbf{w}(t_r)| - \boldsymbol{\beta}(t_r) \cdot (\mathbf{r} - \mathbf{w}(t_r))} \quad (1)$$

$$V(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{q}{|\mathbf{r} - \mathbf{w}(t_r)| - \boldsymbol{\beta}(t_r) \cdot (\mathbf{r} - \mathbf{w}(t_r))} \quad (2)$$

where $\boldsymbol{\beta}(t_r) = \mathbf{v}(t_r)/c$.

$$\mathbf{A}(\mathbf{r}, t) = \frac{\mathbf{v}(t_r)}{c^2} V(\mathbf{r}, t) \quad (3)$$

\mathbf{E} and \mathbf{H} from point charge

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\hat{\mathbf{z}}|}{4\pi\epsilon_0(\hat{\mathbf{z}} \cdot \mathbf{u})^3} [(c^2 - v^2)\mathbf{u} + \hat{\mathbf{z}} \times (\mathbf{u} \times \mathbf{a})]$$

$$\mathbf{B}(\mathbf{r}, t) = \frac{1}{c} \hat{\mathbf{z}} \times \mathbf{E}(\mathbf{r}, t)$$

\mathbf{E} and \mathbf{H} from point charge

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\mathbf{z}|}{4\pi\epsilon_0(\mathbf{z} \cdot \mathbf{u})^3} [(c^2 - v^2)\mathbf{u} + \mathbf{z} \times (\mathbf{u} \times \mathbf{a})]$$

Here \mathbf{z} , \mathbf{u} , \mathbf{v} and \mathbf{a} are evaluated at retarded time t_r .

$$\mathbf{z} = \mathbf{r} - \mathbf{w}(t_r)$$

$$|\mathbf{z}| = |\mathbf{r} - \mathbf{w}(t_r)| = c(t - t_r)$$

$$\mathbf{u} = c\hat{\mathbf{z}}(t_r) - \mathbf{v}(t_r)$$

$$\mathbf{v} = \frac{d\mathbf{w}(t_r)}{dt_r}$$

$$\mathbf{a} = \frac{d\mathbf{v}(t_r)}{dt_r}$$

E and H from point charge with constant velocity

Let $\mathbf{a} = \mathbf{0}$. Then

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\mathbf{z}|}{4\pi\epsilon_0(\mathbf{z} \cdot \mathbf{u})^3}(c^2 - v^2)\mathbf{u}$$

Introduce $\mathbf{R} = \mathbf{r} - \mathbf{v}t_r$.

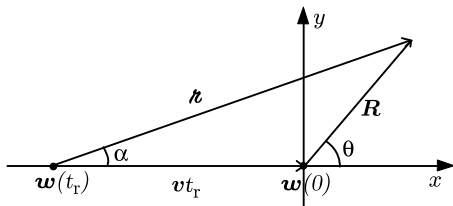


Figure: The particle at $\mathbf{w}(t_r)$ and $\mathbf{w}(0)$.

E and H from point charge with constant velocity

Determine $\mathbf{E}(\mathbf{r}, 0)$ and $\mathbf{B}(\mathbf{r}, 0)$. After some algebra

$$\mathbf{E}(\mathbf{r}, 0) = \frac{q}{4\pi\epsilon_0} \frac{(1 - (v/c)^2)\mathbf{R}}{R^3(1 - v^2 \sin^2 \theta/c^2)^{3/2}} \quad (4)$$

and

$$\mathbf{B}(\mathbf{r}, 0) = \frac{1}{c^2} \mathbf{v} \times \mathbf{E}(\mathbf{r}, 0) \quad (5)$$

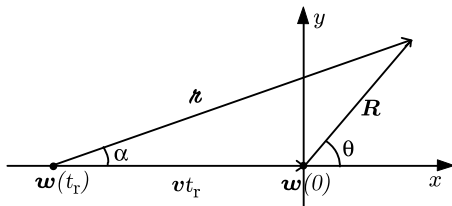


Figure: The particle at $\mathbf{w}(t_r)$ and $\mathbf{w}(0)$.

E and H from point charge with constant velocity

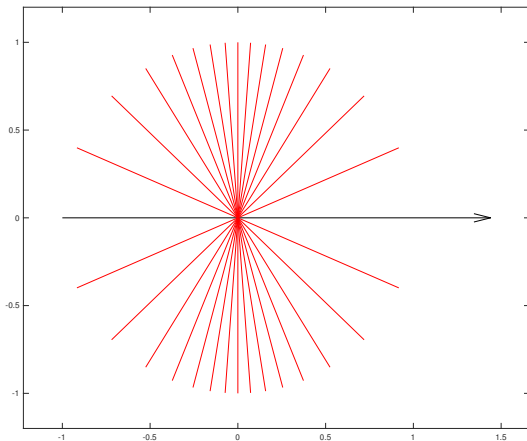


Figure: $E(\mathbf{r}, 0)$ for $\beta = 0.9$.

E and H from point charge with constant velocity

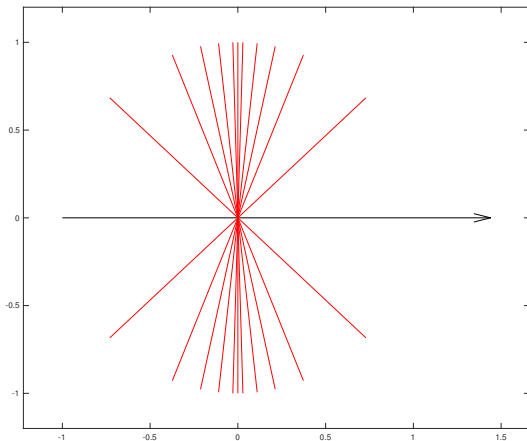


Figure: $E(\mathbf{r}, 0)$ for $\beta = 0.95$.

E and H from point charge with constant velocity

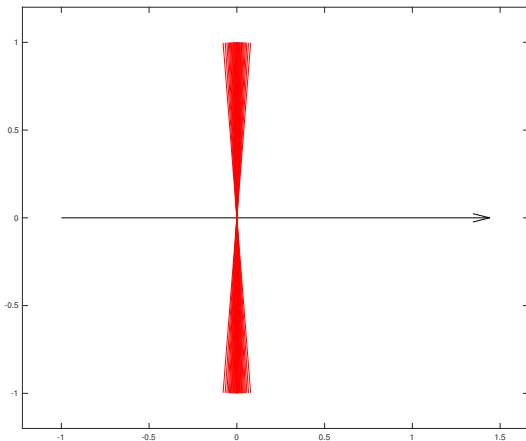


Figure: $E(\mathbf{r}, 0)$ for $\beta = 0.999$.

Today

- ▶ Radiation from currents and charges
- ▶ Larmor formula
- ▶ Liénard's generalization of the Larmor formula

\mathbf{E} and \mathbf{H} from point charge

Question: Can we identify what part of \mathbf{E} that contains the radiated field?

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\mathbf{z}|}{4\pi\epsilon_0(\mathbf{z} \cdot \mathbf{u})^3} [(c^2 - v^2)\mathbf{u} + \mathbf{z} \times (\mathbf{u} \times \mathbf{a})]$$

$$\mathbf{z} = \mathbf{r} - \mathbf{w}(t_r)$$

$$|\mathbf{z}| = |\mathbf{r} - \mathbf{w}(t_r)| = c(t - t_r)$$

$$\mathbf{u} = c\hat{\mathbf{z}}(t_r) - \mathbf{v}(t_r)$$

$$\mathbf{v} = \frac{d\mathbf{w}(t_r)}{dt_r}$$

$$\mathbf{a} = \frac{d\mathbf{v}(t_r)}{dt_r}$$

E and H from point charge

Answer: Yes!

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\mathbf{z}|}{4\pi\epsilon_0(\mathbf{z} \cdot \mathbf{u})^3} \mathbf{z} \times (\mathbf{u} \times \mathbf{a})$$

$$\mathbf{z} = \mathbf{r} - \mathbf{w}(t_r)$$

$$|\mathbf{z}| = |\mathbf{r} - \mathbf{w}(t_r)| = c(t - t_r)$$

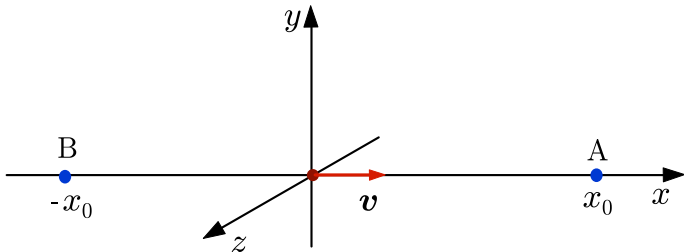
$$\mathbf{u} = c\hat{\mathbf{z}}(t_r) - \mathbf{v}(t_r)$$

$$\mathbf{v} = \frac{d\mathbf{w}(t_r)}{dt_r}$$

$$\mathbf{a} = \frac{d\mathbf{v}(t_r)}{dt_r}$$

E and H from point charge

Assume that $\mathbf{w}(t_r) = \mathbf{0}$ and $\mathbf{v}(t_r) = v\hat{\mathbf{x}}$.



Observer A is at $\mathbf{r}_A = x_0\hat{\mathbf{x}}$, where $x_0 > 0$ and observer B is at $\mathbf{r}_B = -x_0\hat{\mathbf{x}}$.

Who will receive the strongest radiated E -field?

E and H from point charge

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\mathbf{z}|}{4\pi\epsilon_0(\mathbf{z} \cdot \mathbf{u})^3} \mathbf{z} \times (\mathbf{u} \times \mathbf{a})$$

For A

$$\mathbf{z} = x_0 \hat{\mathbf{x}}$$

$$|\mathbf{z}| = x_0$$

$$\mathbf{u} = (c - v) \hat{\mathbf{x}}$$

For B

$$\mathbf{z} = -x_0 \hat{\mathbf{x}}$$

$$|\mathbf{z}| = x_0$$

$$\mathbf{u} = -(c + v) \hat{\mathbf{x}}$$

\mathbf{E} and \mathbf{H} from point charge

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\mathbf{z}|}{4\pi\epsilon_0(\mathbf{z} \cdot \mathbf{u})^3} \mathbf{z} \times (\mathbf{u} \times \mathbf{a})$$

For A

$$|\mathbf{E}(\mathbf{r}, t)| \sim \frac{1}{x_0(c-v)^2}$$

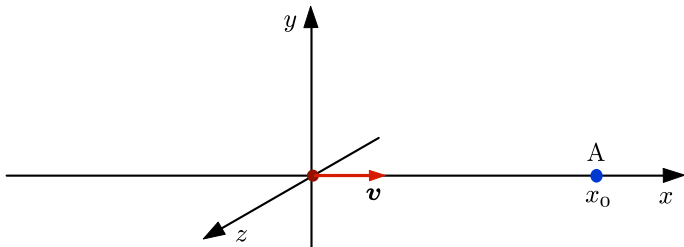
For B

$$|\mathbf{E}(\mathbf{r}, t)| \sim \frac{1}{x_0(c+v)^2}$$

At MAX IV: 3GeV electrons $\Rightarrow \frac{v}{c} = 0.999999986$. Then \mathbf{E} is 10^{16} times larger at A than at B. The power flux is 10^{32} times larger at A than at B.

E and H from point charge

Assume that $\mathbf{w}(t_r) = \mathbf{0}$ and $\mathbf{v}(t_r) = v\hat{\mathbf{x}}$.



Observer A is at $\mathbf{r}_A = x_0\hat{\mathbf{x}}$, where $x_0 > 0$.

In what direction should \mathbf{a} be in order to maximize E at A ?

E and H from point charge

Assume that $\mathbf{w}(t_r) = \mathbf{0}$ and $\mathbf{v}(t_r) = v\hat{\mathbf{x}}$.

Observer A is at $\mathbf{r}_A = x_0\hat{\mathbf{x}}$, where $x_0 > 0$.

In what direction should \mathbf{a} be in order to maximize E at A ?

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\mathbf{z}|}{4\pi\epsilon_0(\mathbf{z} \cdot \mathbf{u})^3} \mathbf{z} \times (\mathbf{u} \times \mathbf{a})$$

Here

$$\mathbf{z} = x_0\hat{\mathbf{x}}$$

$$|\mathbf{z}| = x_0$$

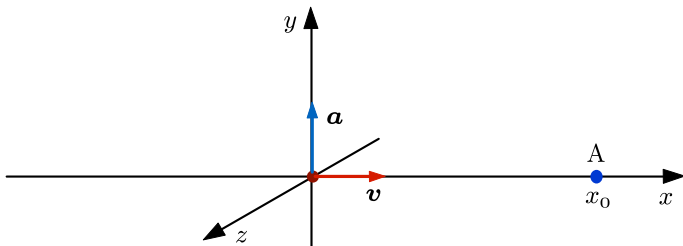
$$\mathbf{u} = (c - v)\hat{\mathbf{x}}$$

$$\mathbf{z} \cdot \mathbf{u} = x_0(c - v)$$

Answer: \mathbf{a} should be perpendicular to $\hat{\mathbf{x}}$.

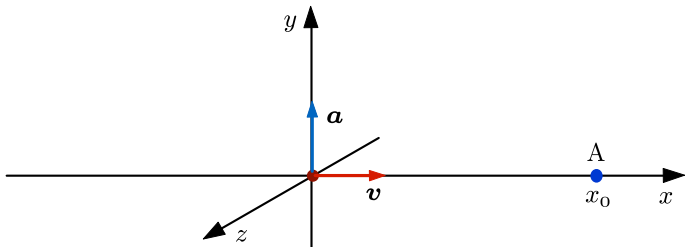
E and H from point charge

$$\mathbf{w}(t_r) = \mathbf{0}, \mathbf{v}(t_r) = v\hat{\mathbf{x}} \text{ and } \mathbf{a}(t_r) = a\hat{\mathbf{y}}$$



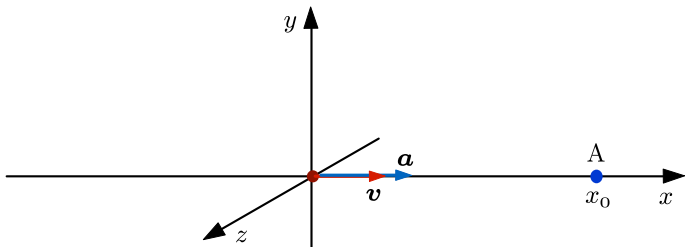
E and H from point charge

The intensity of the light from q is maximized when q travels towards you and its acceleration is perpendicular to the velocity!



E and H from point charge

Question: What is the radiated E at A if $\mathbf{a} = a\hat{x}$?



Answer: Zero

E and H from point charge

Question: What is the radiated \mathbf{E} at A if $\mathbf{a} = a\hat{\mathbf{x}}$?

$$\mathbf{E}(\mathbf{r}, t) = \frac{q|\mathbf{z}|}{4\pi\epsilon_0(\mathbf{z} \cdot \mathbf{u})^3} \mathbf{z} \times (\mathbf{u} \times \mathbf{a})$$

$$\mathbf{z} = x_0\hat{\mathbf{x}}$$

$$|\mathbf{z}| = x_0$$

$$\mathbf{u} = (c - v)\hat{\mathbf{x}}$$

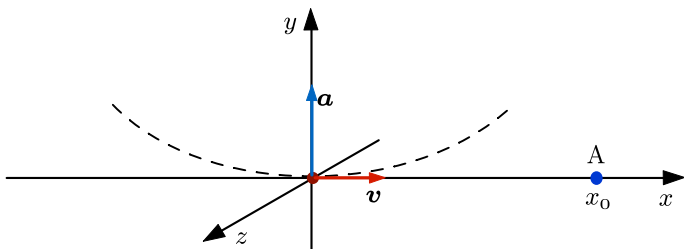
$$\mathbf{v} = v\hat{\mathbf{x}}$$

$$\mathbf{a} = a\hat{\mathbf{x}}$$

Answer: Zero since $\mathbf{u} \times \mathbf{a} = \mathbf{0}$

E and H from point charge

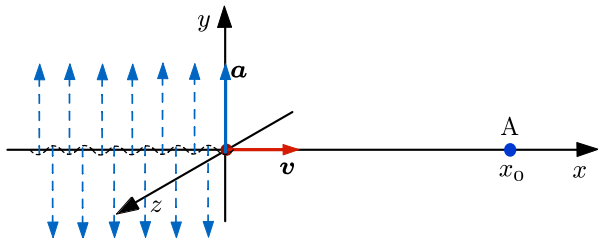
The intensity of the light from q is obtained when q travels towards you and the acceleration is perpendicular to the velocity!



Circular motion in dipole magnet

E and H from point charge

The intensity of the light from q is obtained when q travels towards you and the acceleration is perpendicular to the velocity!



Sinusoidal motion in undulator