

Vector Waves and Probe Compensation

Lecture 9: Probe compensation

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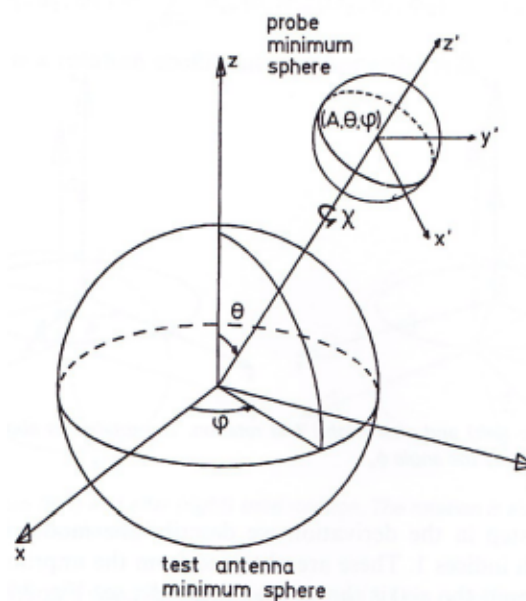
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Experimental setup geometry I

- The two circumscribing (minimum) spheres do not overlap
- Assume there is no coupling between the two antenna systems (no multiple scattering — can be relaxed, see Hansen [2])



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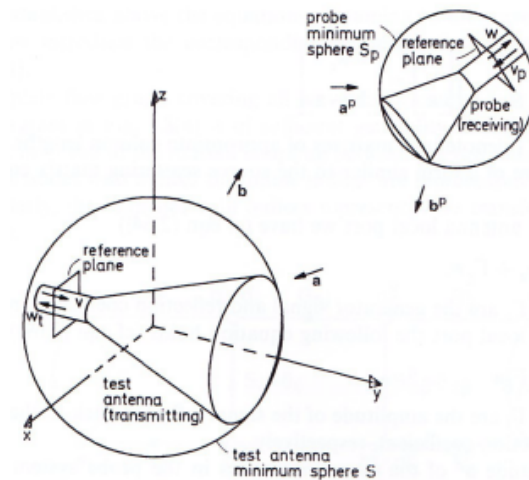
Experimental setup geometry II

Probe scattering description (receiving mode)

$$\begin{pmatrix} w^p \\ b^p \end{pmatrix} = \begin{pmatrix} \Gamma^p & R^p \\ T^p & S^p \end{pmatrix} \begin{pmatrix} v^p \\ a^p \end{pmatrix}$$

Test antenna (AUT) scattering description (transmitting mode)

$$\begin{pmatrix} w^t \\ b^t \end{pmatrix} = \begin{pmatrix} \Gamma^t & R^t \\ T^t & S^t \end{pmatrix} \begin{pmatrix} v_g \\ a^t \end{pmatrix}$$



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Experimental setup geometry III

- The probe is matched to the load, $\Gamma_l = 0$

$$w^p = \frac{1}{1 - \Gamma_l \Gamma^p} \sum_n R_n^p a_n^p = \sum_n R_n^p a_n^p$$

- The AUT are matched to the generator, $\Gamma_g = 0$, and no coupling $a^t = \mathbf{0}$ implies

$$b_n^t = \frac{v_g}{1 - \Gamma_g \Gamma^t} T_n^t = v_g T_n^t$$

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The radiated fields become

$$\begin{cases} \mathbf{E}^t(\mathbf{r}) = k\sqrt{\eta_0\eta}v_g \sum_n T_n^t \mathbf{u}_n(k\mathbf{r}) \\ \mathbf{H}^t(\mathbf{r}) = \frac{k}{i\sqrt{\eta_0\eta}}v_g \sum_n T_n^t \mathbf{u}_{\bar{n}}(k\mathbf{r}) \end{cases} \quad r > r_0$$

in the unprimed coordinate system (r, θ, ϕ) (only sum over $\tau = 1, 2$)

The truncation size is

$$J = 2L(L + 2) \text{ where } l_{\max} = L = [kr_0] + n_1 \text{ (empirical)}$$

where r_0 is the radius of the circumscribing sphere, and $[x]$ is the integer part of x

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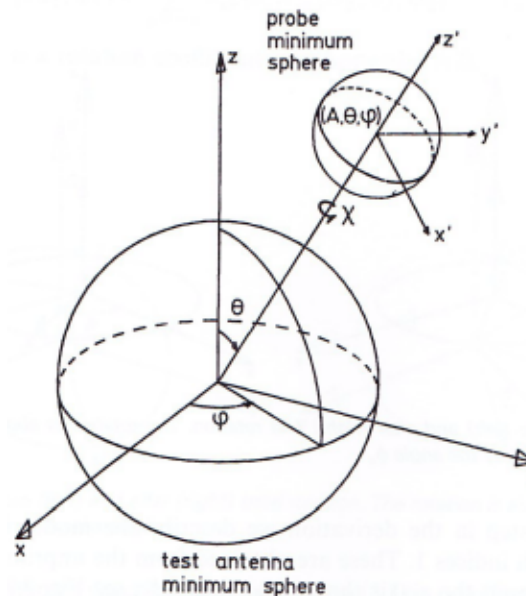
Goal: To find the transmission coefficients T_n^t of the test antenna

Aim: Transfer this description to the coordinate system of the probe!

The transformation consists of three parts, **1) rotation**, **2) translation along z axis**, and a final **3) rotation**

Transmitting AUT, receiving probe III

- 1 Rotations ϕ and θ
- 2 Translation A along the new z axis
- 3 Final rotation χ around the z axis



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- 1 **Rotation:** From Lectures 5 & 6 we have

$$\begin{Bmatrix} \mathbf{v}_n(k\mathbf{r}) \\ \mathbf{u}_n(k\mathbf{r}) \end{Bmatrix} = \sum_{n'} \mathcal{D}_{n'n}(\alpha, \beta, \gamma) \begin{Bmatrix} \mathbf{v}'_{n'}(k\mathbf{r}') \\ \mathbf{u}'_{n'}(k\mathbf{r}') \end{Bmatrix}$$

$$\mathcal{D}_{n'n}(\alpha, \beta, \gamma) = \delta_{\tau\tau'} \delta_{ll'} \delta_{mm'} e^{im'\gamma} d_{m'm}^{(l)}(\beta) e^{im\alpha}$$

$$d_{m'm}^{(l)}(\beta) = \sqrt{\frac{(l+m')!(l-m')!}{(l+m)!(l-m)!}} \sum_j (-1)^{l-m'-j}$$

$$\binom{l+m}{l-m'-j} \binom{l-m}{j} \left(\cos \frac{\beta}{2}\right)^{2j+m+m'} \left(\sin \frac{\beta}{2}\right)^{2l-2j-m-m'}$$

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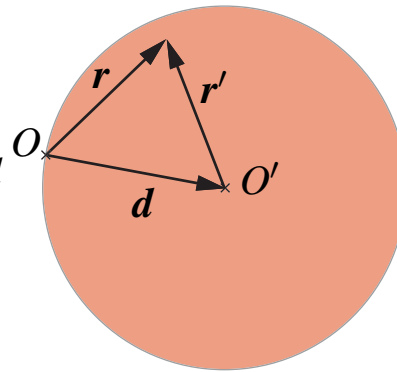
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2 Translation:

$$\mathbf{u}_n(k\mathbf{r}) = \sum_{n'} \mathcal{P}_{nn'}(k\mathbf{d}) \mathbf{v}_{n'}(k\mathbf{r}'), \quad r' < d$$



The matrix $\mathcal{P}_{nn'}(k\mathbf{d})$ is given by

$\tau \backslash \tau'$	1	2	3
1	$A_{nn'}(k\mathbf{d}; h_\lambda^{(1)})$	$B_{nn'}(k\mathbf{d}; h_\lambda^{(1)})$	0
2	$B_{nn'}(k\mathbf{d}; h_\lambda^{(1)})$	$A_{nn'}(k\mathbf{d}; h_\lambda^{(1)})$	0
3	0	0	$C_{nn'}(k\mathbf{d}; h_\lambda^{(1)})$

$A_{nn'}(k\mathbf{d}; h_\lambda^{(1)})$, $B_{nn'}(k\mathbf{d}; h_\lambda^{(1)})$, and $C_{nn'}(k\mathbf{d}; h_\lambda^{(1)})$ are explicitly given in Lecture 6

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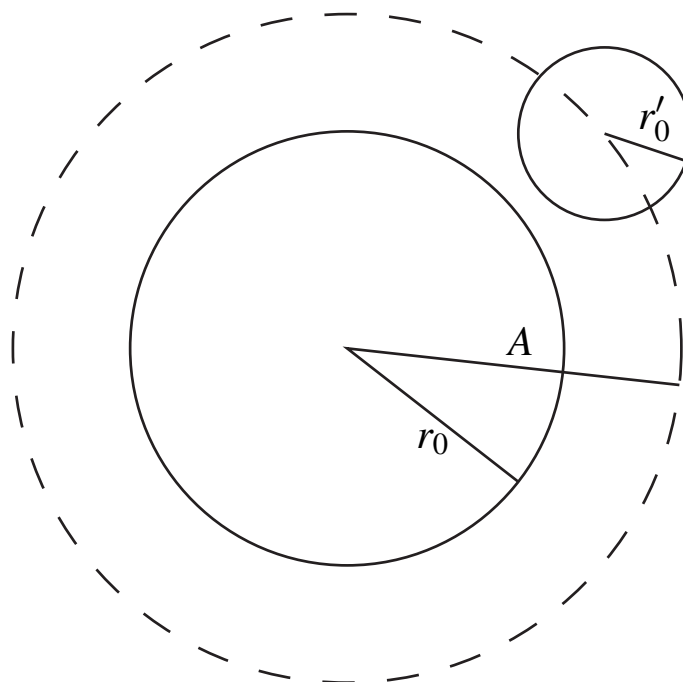
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The connection between the two coordinate systems (r, θ, ϕ) and (r', θ', ϕ') are written in a compact notation

$$\mathbf{u}_n(k\mathbf{r}) = \sum_{n'} \mathcal{C}_{nn'}(k\mathbf{d}) \mathbf{v}'_{n'}(k\mathbf{r}'), \quad r' < (A - r_0)$$

where $\mathcal{C}_{nn'}(k\mathbf{d})$ is a combination of matrices that represent the appropriate rotations and translation

$$\begin{aligned} \mathcal{C}_{nn'}(k\mathbf{d}) &= \sum_{n''n'''} \mathcal{D}_{n''n}(\phi, \theta, 0) \mathcal{P}_{n''n'''}(kA\hat{\mathbf{z}}) \mathcal{D}_{n'n'''}(0, 0, \chi) \\ &= e^{im\phi} d_{m'm}^{(l)}(\theta) \mathcal{P}_{\tau lm' \tau' l' m'}(kA\hat{\mathbf{z}}) e^{im'\chi} \end{aligned}$$

Remember $\mathcal{D}_{n'n'''}(0, 0, \chi)$ and $\mathcal{D}_{n''n}(\phi, \theta, 0)$ diagonal in τ and l indices, and $\mathcal{P}_{n''n'''}(kA\hat{\mathbf{z}})$ diagonal in the m index

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$$\mathbf{E}^t(\mathbf{r}) = k\sqrt{\eta_0\eta}v_g \sum_n T_n^t \mathbf{u}_n(k\mathbf{r}) = k\sqrt{\eta_0\eta}v_g \sum_{nn'} T_n^t \mathcal{C}_{nn'}(k\mathbf{d}) \mathbf{v}'_{n'}(k\mathbf{r}')$$

This is the description of the electric field in the primed system (r', θ', ϕ')

From $2\mathbf{v}'_n(k\mathbf{r}') = \mathbf{u}'_n(k\mathbf{r}') + \mathbf{w}'_n(k\mathbf{r}')$, we get $(r' < (A - r_0))$

$$\mathbf{E}^t(\mathbf{r}) = \frac{k\sqrt{\eta_0\eta}v_g}{2} \sum_{nn'} T_n^t \mathcal{C}_{nn'}(k\mathbf{d}) (\mathbf{u}'_n(k\mathbf{r}') + \mathbf{w}'_n(k\mathbf{r}'))$$

Identify the input, $a_{n'}^p$, in the primed system

$$a_{n'}^p = \frac{v_g}{2} \sum_n T_n^t \mathcal{C}_{nn'}(k\mathbf{d})$$

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In the primed system, the antenna scattering description is (matched conditions $\Gamma^p = 0$)

$$w^p(A, \chi, \theta, \phi) = \sum_{n'} R_{n'}^p a_{n'}^p = \frac{V_g}{2} \sum_{nn'} T_n^t C_{nn'}(kd) R_{n'}^p$$

where R_n^p are the receiving coefficients of the probe, which are assumed known from a previous calibration process

Problem in a nutshell: From the knowledge of the received wave $w^p(A, \chi, \theta, \phi)$, determine the unknown transmitting coefficients T_n^t

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 - Analytical methods
 - Full field
 - Radial field component [2]
 - Wood's method [2]
- With probe compensation
 - Naïve methods
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Without probe compensation

These methods rely on the assumption that the electric and the magnetic fields are accurately measured with e.g., short electric and magnetic dipoles

$$\begin{cases} \mathbf{E}^t(\mathbf{r}) = k\sqrt{\eta_0\eta} \sum_n b_n^t \mathbf{u}_n(k\mathbf{r}) \\ \mathbf{H}^t(\mathbf{r}) = \frac{k}{i\sqrt{\eta_0\eta}} \sum_n b_n^t \mathbf{u}_{\bar{n}}(k\mathbf{r}) \end{cases} \quad r > r_0$$

The probe should have **low** directivity and located at not-to-close distance

Goal: Find the expansion coefficients b_n^t

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Naïve approach

Measure the fields, $\mathbf{E}^t(\mathbf{r}_i)$, at a set of points \mathbf{r}_i , $i = 1, 2, \dots, N$ ($\hat{\mathbf{e}}_i$ given unit vector, may vary with the points \mathbf{r}_i)

$$\hat{\mathbf{e}}_i \cdot \mathbf{E}^t(\mathbf{r}_i) = k\sqrt{\eta_0\eta} \sum_n b_n^t \hat{\mathbf{e}}_i \cdot \mathbf{u}_n(k\mathbf{r}_i) = \sum_n P_{in} b_n^t, \quad i = 1, \dots, N$$

where the known matrix P_{in} is

$$P_{in} = k\sqrt{\eta_0\eta} \hat{\mathbf{e}}_i \cdot \mathbf{u}_n(k\mathbf{r}_i)$$

Procedure: Measure a set where $N \geq J$ (J is the truncation level of spherical vector waves) and solve the overdetermined system with SVD

Hansen [2] is sceptical to this method — depends critically on the sampling

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Assume accurate measurements of the electric field on a spherical surface $r = \text{constant} > r_0$ are available

$$\begin{cases} \mathbf{E}^t(\mathbf{r}) = k\sqrt{\eta_0\eta} \sum_n b_n^t \mathbf{u}_n(k\mathbf{r}) \\ \mathbf{H}^t(\mathbf{r}) = \frac{k}{i\sqrt{\eta_0\eta}} \sum_n b_n^t \mathbf{u}_{\bar{n}}(k\mathbf{r}) \end{cases} \quad r > r_0$$

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Scalar multiplication with $\mathbf{A}_{1n}^\dagger(\hat{\mathbf{r}})$ and integrate over the unit sphere Ω (in practice numerical integration)

$$\begin{cases} \iint_{\Omega} \mathbf{A}_{1n}^\dagger(\hat{\mathbf{r}}) \cdot \mathbf{E}^t(r\hat{\mathbf{r}}) \, d\Omega = k\sqrt{\eta_0\eta} b_{1n}^t h_l^{(1)}(kr) \\ \iint_{\Omega} \mathbf{A}_{2n}^\dagger(\hat{\mathbf{r}}) \cdot \mathbf{E}^t(r\hat{\mathbf{r}}) \, d\Omega = k\sqrt{\eta_0\eta} b_{2n}^t \frac{(krh_l^{(1)}(kr))'}{kr} \end{cases} \quad r > r_0$$

Solve for b_n ($h_l^{(1)}(z)$ and $(krh_l^{(1)}(kr))'$ have no real zeros, see Kristensson [4, Lemma B.1 & B.2])

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Hansen [2] multiplies with $\mathbf{u}_n^\dagger(kr\hat{\mathbf{r}})$ and integrates over the unit sphere Ω

The result is

$$\left\{ \begin{aligned} \iint_{\Omega} \mathbf{u}_{1n}^\dagger(kr\hat{\mathbf{r}}) \cdot \mathbf{E}^t(r\hat{\mathbf{r}}) \, d\Omega &= k\sqrt{\eta_0\eta}b_{1n}^t \left(h_l^{(1)}(kr) \right)^2 \\ \iint_{\Omega} \mathbf{u}_{2n}^\dagger(kr\hat{\mathbf{r}}) \cdot \mathbf{E}^t(r\hat{\mathbf{r}}) \, d\Omega &= k\sqrt{\eta_0\eta}b_{2n}^t \left\{ \left(\frac{(krh_l^{(1)}(kr))'}{kr} \right)^2 \right. \\ &\quad \left. + l(l+1) \left(\frac{h_l^{(1)}(kr)}{kr} \right)^2 \right\} \end{aligned} \right. \quad r > r_0$$

These two approaches are very similar

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From above

$$w^p(A, \chi, \theta, \phi) = \frac{v_g}{2} \sum_{nn'} T_n^t C_{nn'}(k\mathbf{d}) R_n^p$$

where

$$C_{nn'}(k\mathbf{d}) = e^{im\phi} d_{m'm}^{(l)}(\theta) \mathcal{P}_{\tau lm' \tau' l' m'}(kA\hat{\mathbf{z}}) e^{im'\chi}$$

Notice that the sensitivity to different polarization of the probe signal is included in the receiving coefficients R_n^p

As above, the indices are:

$$\tau = 1, 2, \quad l = 1, 2, \dots, L, \quad m = -l, \dots, 0, \dots, l$$

$$\tau' = 1, 2, \quad l' = 1, 2, \dots, L', \quad m' = -l', \dots, 0, \dots, l'$$

Rule of thumb Hansen [2]: $L = [kr_0] + 10$ and $L' = [kr'_0] + 10$

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Naïve approach

Measure the response, $w^p(A_i, \chi_i, \theta_i, \phi_i)$, at a set of points r_i or $(A_i, \chi_i, \theta_i, \phi_i)$ (the method allows for different A_i), $i = 1, 2, \dots, N$

$$w^p(A_i, \chi_i, \theta_i, \phi_i) = \frac{v_g}{2} \sum_{nn'} T_n^t C_{nn'}(kr_i) R_{n'}^p = \sum_n P_{in} T_n^t$$

where the known matrix P_{in} is

$$P_{in} = \frac{v_g}{2} \sum_{n'} C_{nn'}(kr_i) R_{n'}^p$$

Procedure: Measure a set where $N \geq J$ (J is the truncation level of spherical vector waves) and solve the overdetermined system with SVD

Hansen [2] claims this method is ill-conditioned — depends critically on the sampling

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Rewrite as

$$\begin{aligned} w^p(A, \chi, \theta, \phi) &= \frac{v_g}{2} \sum_{nn'} T_n^t C_{nn'}(kd) R_{n'}^p \\ &= \frac{v_g}{2} \sum_{nn'} T_n^t e^{im\phi} d_{m'm}^{(l)}(\theta) \mathcal{P}_{\tau lm' \tau' l' m'}(kA \hat{z}) e^{im' \chi} R_{n'}^p \\ &= v_g \sum_{\tau l m m'} T_{\tau l m}^t e^{im\phi} d_{m'm}^{(l)}(\theta) e^{im' \chi} P_{\tau l m'}(kA) \end{aligned}$$

where the **probe response constants** are

$$P_{\tau l m'}(kA) = \frac{1}{2} \sum_{\tau' l'} \mathcal{P}_{\tau l m' \tau' l' m'}(kA \hat{z}) R_{\tau' l' m'}^p, \quad m' = -L', \dots, L'$$

These constants contain only known quantities — specific for each range and probe

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We assume measurement are made in a near-field range on a spherical surface $A = \text{constant}$

Three types of integrations retrieve the transmitting coefficients $T_{\tau lm}^t$

- 1 Integrate over χ over $[0, 2\pi)$ (polarization)
- 2 Integrate over the azimuth angle ϕ over $[0, 2\pi)$
- 3 Integrate over the polar angle θ over $[0, \pi]$

These three integrations are done numerical by proper sampling, see Hansen [2]

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1. First step (the polarization χ)

With the use of orthogonality of $e^{im'\chi}$ over $[0, 2\pi)$

$$w_{m'}^p(A, \theta, \phi) = v_g \sum_{\tau lm} T_{\tau lm}^t e^{im\phi} d_{m'm}^{(l)}(\theta) P_{\tau lm'}(kA)$$

where

$$w_{m'}^p(A, \theta, \phi) = \frac{1}{2\pi} \int_0^{2\pi} w^p(A, \chi, \theta, \phi) e^{-im'\chi} d\chi$$

and (approximation)

$$w^p(A, \chi, \theta, \phi) = \sum_{m'=-L'}^{L'} w_{m'}^p(A, \theta, \phi) e^{im'\chi}$$

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2. Second step (the azimuth angle ϕ)

With the use of orthogonality of $e^{im\phi}$ over $[0, 2\pi)$

$$w_{mm'}^p(A, \theta) = v_g \sum_{\tau l} T_{\tau lm}^t d_{m'm}^{(l)}(\theta) P_{\tau lm'}(kA)$$

where

$$w_{mm'}^p(A, \theta) = \frac{1}{2\pi} \int_0^{2\pi} w_{m'm}^p(A, \theta, \phi) e^{-im\phi} d\phi$$

and (approximation)

$$w^p(A, \chi, \theta, \phi) = \sum_{m=-L}^L \sum_{m'=-L'}^{L'} w_{mm'}^p(A, \theta) e^{im'\chi} e^{im\phi}$$

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3. Third step (the polar angle θ)

With the use of orthogonality of $d_{m'm}^{(l)}(\theta)$ over $[0, \pi]$

$$w_{lmm'}^p(A) = v_g \sum_{\tau} T_{\tau lm}^t P_{\tau lm'}(kA) \quad (\star)$$

where

$$w_{lmm'}^p(A) = \frac{2l+1}{2} \int_0^{\pi} w_{mm'}^p(A, \theta) d_{m'm}^{(l)}(\theta) \sin \theta d\theta$$

We have here used the orthogonality of Wigner's function $d_{m'm}^{(l)}(\theta)$, see Appendix and Edmonds [1]

$$\int_0^{\pi} d_{m'm}^{(l)}(\theta) d_{m'm}^{(l)}(\theta) \sin \theta d\theta = \delta_{ll'} \frac{2}{2l+1}$$

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Final, approximate expression

$$w^p(A, \chi, \theta, \phi) = \sum_{l=0}^L \sum_{m=-L}^L \sum_{m'=-L'}^{L'} w_{lmm'}^p(A) e^{im'\chi} e^{im\phi} d_{m'm}^{(l)}(\theta)$$

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4. Fourth step (the transmitting coefficients T_n^t)

Let $m' = \pm 1$ in (*), Wacker [6], (first-order probe compensation) *i.e.*,

$$\begin{cases} v_g T_{1lm}^t P_{1l-1}(kA) + v_g T_{2lm}^t P_{2l-1}(kA) = w_{lm-1}^p(A) \\ v_g T_{1lm}^t P_{1l1}(kA) + v_g T_{2lm}^t P_{2l1}(kA) = w_{lm1}^p(A) \end{cases}$$

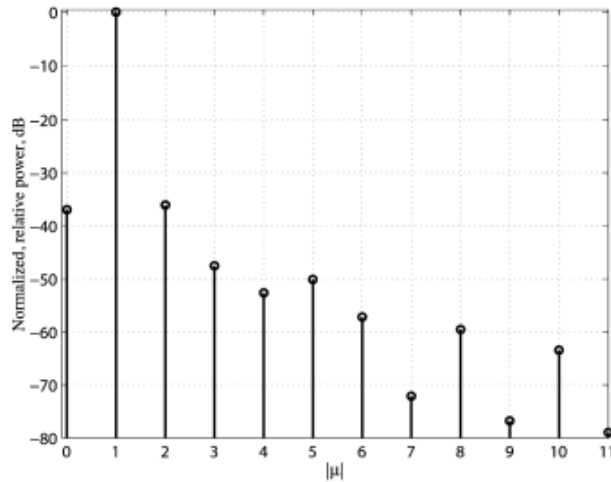
Solve this 2×2 system for each l, m values for the two unknowns T_{1lm}^t and T_{2lm}^t , v_g known

This can be done provided

$$\begin{vmatrix} P_{1l-1}(kA) & P_{2l-1}(kA) \\ P_{1l1}(kA) & P_{2l1}(kA) \end{vmatrix} \neq 0$$

Higher-order probe compensation I

In practice, real probes are never ideal first-order probes but possess, due to, e.g., the manufacturing errors, also a minor level of high-order modes [5]



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Higher-order probe compensation II

This requires higher-order probe compensation (hopefully more at DTU)

We follow [5] and write (superscripts p and t are omitted)

$$w(A, \chi, \theta, \phi) = v_g \sum_{\tau l m m'} T_{\tau l m} e^{im\phi} d_{m'm}^{(l)}(\theta) e^{im'\chi} P_{\tau l m'}(kA)$$

as $w(A, \chi, \theta, \phi) = w_d(A, \chi, \theta, \phi) + w_e(A, \chi, \theta, \phi)$

$$w_d(A, \chi, \theta, \phi) = v_g \sum_{\tau l m} \sum_{m'=\pm 1} T_{\tau l m} e^{im\phi} d_{m'm}^{(l)}(\theta) e^{im'\chi} P_{\tau l m'}(kA) = \Phi_d\{T\}$$

and

$$w_e(A, \chi, \theta, \phi) = v_g \sum_{\tau l m} \sum_{m' \neq \pm 1} T_{\tau l m} e^{im\phi} d_{m'm}^{(l)}(\theta) e^{im'\chi} P_{\tau l m'}(kA) = \Phi_e\{T\}$$

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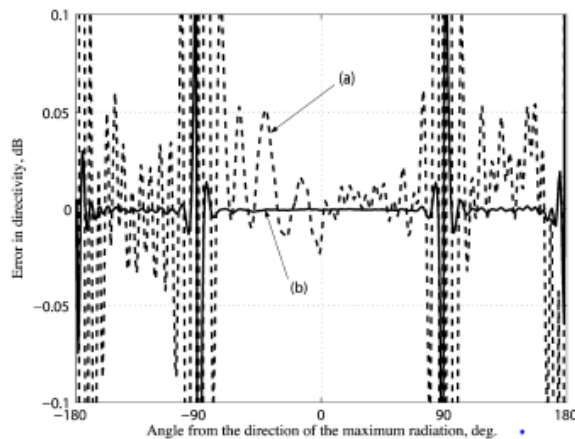


Higher-order probe compensation III

Formally, create an iteration scheme

$$\begin{cases} T^{(i)} = \Phi_d^{-1}\{w - w_e^{(i)}\} \text{ (first-order probe compensation)} \\ w_e^{(i+1)} = \Phi_e\{T^{(i)}\} \end{cases}$$

The iteration index $i = 0, 1, \dots, I$, and the iteration scheme is initialized by $w_e^{(0)} = 0$



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What has not been covered, Hansen [2]

- Friis' transmission formula
- Far field aspects and transformations
- Measurement techniques
- Experimental setups
- Discretization and sampling
- Multiple coupling
- More on higher-order probe compensation, maybe solved at DTU
- ...
- ...
- ...

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Appendix, Wigner's function $d_{mm'}^{(l)}(\theta)$

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Rewrite in terms of Jacobi polynomials $P_m^{(\alpha,\beta)}(\cos \theta)$

$$d_{m'm}^{(l)}(\theta) = \sqrt{\frac{(l+m')!(l-m')!}{(l+m)!(l-m)!}} \left(\cos \frac{\theta}{2}\right)^{m'+m} \left(\sin \frac{\theta}{2}\right)^{m'-m} P_{l-m'}^{(m'-m, m'+m)}(\cos \theta)$$

and use the orthogonality relation of these polynomials [3, Lemma 5.8]

$$\int_{-1}^1 (1-x)^\alpha (1+x)^\beta P_n^{(\alpha,\beta)}(x) P_m^{(\alpha,\beta)}(x) dx = \frac{\delta_{n,m} 2^{\alpha+\beta+1} \Gamma(\alpha+n+1) \Gamma(\beta+n+1)}{n! (\alpha+\beta+2n+1) \Gamma(\alpha+\beta+n+1)}$$

$$\begin{aligned} \int_0^\pi d_{m'm}^{(l)}(\theta) d_{m'm}^{(l')}(\theta) \sin \theta d\theta &= \sqrt{\frac{(l+m')!(l-m')!}{(l+m)!(l-m)!}} \sqrt{\frac{(l'+m')!(l'-m')!}{(l'+m)!(l'-m)!}} \\ &\times \int_{-1}^1 \left(\frac{1-x}{2}\right)^{m'-m} \left(\frac{1+x}{2}\right)^{m'+m} P_{l-m'}^{(m'-m, m'+m)}(x) P_{l'-m'}^{(m'-m, m'+m)}(x) dx \\ &= \delta_{l'l} 2 \frac{(l+m')!(l-m')!}{(l+m)!(l-m)!} \frac{(l-m)!(l+m)!}{(l-m')!(2l+1)(l+m')!} = \delta_{l'l} \frac{2}{2l+1} \end{aligned}$$

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Make a probe compensation on a electric dipole system — both the AUT and the probe are electric dipoles

Perform the following steps:

- 1 Determine the transmitting and receiving coefficients of an electric dipole oriented along the three Cartesian axes (use the result of a vertical dipole and the result of the Assignment 5 to derive the two remaining orientations, *cf.*, with the results in Hansen [2])
- 2 Determine the received signal, $w^p(A, \chi, \theta, \phi)$, from a vertical electrical dipole in spherical near field range if the probe is an electrical dipole (do not expand the translation coefficients more than necessary to perform the next step)
- 3 From the result in item 2, make a probe compensation to retrieve the original AUT coefficients